Single-Molecule Techniques in Biological and Biophysical Research Gerhard A. Bab

# Single-Molecule Techniques in Biological and Biophysical Research

# Proefschrift

ter verkrijging van de graad van Doctor aan de Universiteit Leiden, op gezag van de Rector Magnificus Dr. D. D. Breimer, hoogleraar in de faculteit der Wiskunde en Natuurwetenschappen en die der Geneeskunde, volgens besluit van het College voor Promoties te verdedigen op donderdag 8 januari 2004 klokke 14.15 uur

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# Dedication

Dedicated to the memory of Professor Dr. Hansgeorg Schindler (\*September 27, 1943 – †August 28, 2001)

Möge er Antworten gefunden haben auf all seine "kurzen Fragen".



#### DEDICATION

Professor Dr. Hansgeorg Schindler war der Gründer des Instituts für Biophysik an der Johannes Kepler Universität Linz, Österreich, und mein erster Lehrer auf diesem Gebiet.

Sein Markenzeichen, ein kurz vor dem Mittagessen geäußertes "Könnten Sie eben in mein Büro kommen, ich hätte eben eine kurze Frage!", war der übliche Auftakt für lange Gespräche über den Fortgang der Arbeit im Labor. Er war sich auch der wirtschaftlichen Aspekte der Wissenschaft bewusst, und mit Hilfe des Landes Oberösterreich gründete er die "Upper Austrian Research", einen Biotech-Betrieb zur Entwicklung und Vermarktung ultra-sensitiver Mikroskopiemethoden.

Am 28. August 2001 verstarb Professor Schindler an den Folgen eines Freizeitunfalls während eines Urlaubs in Frankreich.

Professor Dr. Hansgeorg Schindler was the founder of the Biophysics institute at the Johannes Kepler University in Linz, Austria, as well as my first teacher in this field.

His trademark comment "Could you come to my office for a moment, I've just got a short question!", uttered shortly before lunch, usually meant a long scientific discussion about the progress in the lab. He was also well aware of the commerial side of science. With financial help from the state of Upper Austria he founded the biotech company "Upper Austrian Research", whose aim is to develop and market ultra-sensitive microscopy methods.

Professor Schindler died on August 28, 2001, in an accident occuring during a vacation in France.

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Alle Kriegsherren haben einen gemeinsamen Feind: die Wahrheit.

Kein Volk ist besser oder schlechter als dein eigenes.

Jeder Krieg ist eine Niederlage. Denn Krieg vernichtet Leben.

Wer Kriege im Namen Gottes führt, ist stets des Teufels.

Es gibt weder gerechte noch heilige Kriege.

All warlords have one common enemy: the Truth.

No people are better or worse than your own.

The outcome of every war is a defeat, because war destroys life.

Anyone who wages war in the name of God will always end up in hell.

There are no such things as Just or Holy Wars.

Kurt Tucholsky (1890 – 1935) "Wenn sie wieder lügen"

# Chapter 1

# Introduction

Long before Richard P. Feynman brought the world of nanostructures into the scientific spotlight by coining the phrase "There's plenty of room at the bottom," researchers have tried to elucidate with varying degrees of success what this "bottom" actually looks like.

Since the early days of microscopy, when microscopes consisted of relatively simple lens systems, the aim was to find more detail, reach higher resolution, as every new insight gained by resolving smaller detail will immediately lead to new questions. When Anton van Leeuwenhoek discovered the first protozoa in 1674 or Schleiden and Schwann proposed cells as building blocks of animals (1838), the question arose what these protozoa and cells made were of.

In the search for images with greater and greater detail, scientist were continually forced to overcome barriers set by their equipment. When the technical circumstances were right and optical systems could be manufactured reproducibly with great accuracy, the barrier shifted toward limits set by the wave nature of light. These resolution limits are determined by the wavelength,  $\lambda$ , of light, giving access to length scales on the order of  $\lambda/2$ .

A more accurate value for the resolution is  $1.22 \frac{\lambda}{2 \text{ NA}}$ , derived from the Fraunhofer diffraction of a circular aperture for a given wavelength  $\lambda$  and a numerical aperture of the objective NA. The theoretical discussion of the resolving power of a microscope by Abbé [1] dates back to the late  $19^{th}$  century. A detailed overview of microscopy methods and basics can be found in Pawlay [2] or Abramowitz [3]. A historical overview is given in Alberts et al. [4].

## **1.1** Advancements in Microscopy

#### 1.1.1 From in vitro...

	wide-field	$\operatorname{confocal}$
lateral	$0.52 \ \frac{\lambda}{NA}$	$0.37 \ \frac{\lambda}{NA}$
axial	1.77 $\frac{\lambda}{NA^2}$	1.28 $\frac{\lambda}{NA^2}$

Table 1.1: The Full-Width-Half-Maximum (FWHM) of the point-spread functions of optical wide-field microscopy and confocal microscopy. Technical advances have allowed to shift this barrier set by the wavelength of light further to smaller and smaller structures, by using shorter and shorter wavelengths in the ultra-violet part of the spectrum, or by applying new focusing techniques such as confocal[2], and 4Pi-confocal[5]. Alternatively one can also employ multi-photon processes[6, 7], Point-Spread Function (PSF)-

engineering[8] or computational methods to evaluate finer details of a specimen. A way to truly break the wavelenght-limit is the use of near-field methods[9].

Consequently one can leave the domain of classical optics altogether and use an electron microscope, gaining the power to resolve single atoms in a lattice.

The price for these advantages, however, was to abandon biologically benign sample conditions and use harmful radiation (UV) or measure in vacuum (EM). These treatments of biological material disturb biological processes heavily, rendering most measurements to actually being done *in mortui*.

#### 1.1.2 ... to in vivo

New methods had to be developed to overcome the need for destructive preparations and return to more biologically relevant conditions. A viable means to do this was to literally feel instead of look. Atomic Force Microscopy (AFM) uses a sharpened needle with a tip radius of only a few nanometers and a highly advanced force feedback to "feel" with resolutions down to sub-nanometers. Even under non-ideal circumstances this method gives access to proteins and even subunits thereof.

However, this method has its drawbacks as it is restricted to the surface of a given sample. On the other hand, the tactile approach allows to gain more than just topographical information. It will also reveal details about specific and unspecific interactions between the sample surface and the clean or coated tip, and measure the elastic properties of the sample. In the case of STM[10], which works on similar methods, yet with a current feedback to ascertain the topography, it is even possible to probe the electrical properties of e.g. ion channels or ion pumps in biological



membranes[11].

Single molecule fluorescence microscopy<sup>1</sup> is an alternative way to gain positional information below the resolution barrier set by the wavelength of light. The first report of the detection of a single biological molecule, an antibody tagged with about 100 fluorophores[15, 16] dates to 1976. Since then the sensitivity of the detection apparatus has been refined so that today the detection of the fluorescence emitted by a single fluorophore can be localized and quantified[17]

As a single fluorophore is essentially a point source on the length scale of resolution for a light microscopy, one can accurately pinpoint the center of the radial point-spread function, h(u, v)

$$h(u,v) = I(v) \equiv (\pi N) \left(\frac{2J_1(v)}{v}\right)^p \tag{1.1}$$

where  $v = \frac{2\pi}{\lambda} r \cdot \text{NA}$  is the optical coordinate in radial, and  $u = \frac{2\pi}{\lambda} z \frac{\text{NA}^2}{n}$ axial direction,  $N = \frac{a^2}{\lambda f}$  the Fresnel number,  $\lambda$  the wavelength of of light used, and a, f, NA are the aperture size, focal length and numerical aperture of the objective.  $J_1(v)$  is the Bessel function of the first kind. The power p in equation 1.1 is determined by the type of microscopy used, i.e. 2 for conventional and 4 for confocal, in effect giving the latter a 30% better resolution.

As mentioned before, in the case of a single emitting fluorophore - or several fluorophores confined to a space much smaller than the resolution mentioned above - one can localize the position of this fluorescent marker with an accuracy much higher than the actual resolution of the microscope[18].

<sup>&</sup>lt;sup>1</sup>for a general overview over the field of single molecule fluorescence microscopy, please refer to the reviews by Moerner and Orrit [12], Weiss [13], and Sako and Yanagida [14].

A closed form for the estimation of the localization accuracy is presented by Thompson et al. [19], taking into account the spatial sampling of the detector (pixel size, a), the background noise b, the standard deviation of the point-spread function s and the number of collected photons N:

$$\left\langle (\Delta x)^2 \right\rangle = \frac{s^2 + \frac{a^2}{12}}{N} + \frac{8\pi s^4 b^2}{a^2 N^2}.$$
 (1.2)

For the typical background in *in vivo* systems and the technical parameters of our imaging system, we can reach an accuracy of about 40 nm. Recent studies by Yildiz et al. [20] even show that *in vitro* an accuracy of down to 1.5 nm can be achieved.

In recent years the focus of imaging biophysics has shifted from "how" to "what". The techniques have grown ever more sophisticated and robust, the ultimate sensitivity of single molecules has been obtained, the know-"how" is there. Now these methods have to be applied to biological research.

## **1.2** Theory of Single Molecule Fluorescence

A fluorophore can generally be described in terms of a four-level system as schematically shown in 1.1. The intensity dependent excitation rate  $k_{exc} = \sigma(\lambda)I$  raises electrons from the ground state  $|G\rangle$  to a vibrational level in the excited state  $|E\rangle$ . A fast relaxation ( $\leq 1$  ps) leaves the electron in the lowest excited singlet level, from which it can return to the ground state with a rate of  $k_s = 1/\tau_s$  or switch over to a triplet state  $|T\rangle$  in a process designated "Inter System Crossing" (ISC). This process is spin forbidden, as both ground and excited state have spin S = 0, while the triplet has



S = 1, but can occur by mediation of spin-orbit coupling.<sup>2</sup>

As the return from the triplet to the ground state is obviously also spinforbidden, it is usually several orders of magnitude slower than the decay from the excited state. If a photon is emitted from the triplet state, the process is called phosphorescence. In principle, both fluorescence and phosphorescence can occur in the same molecule, but due to the slow rate of phosphorescence, non-ratiative relaxation from the triplet to the ground state is often the dominant process.



Figure 1.1: A simplified Jabłoński diagram of the main energetic levels in a fluorophore. Typical values for the rate constants can be found in table 1.2.

Gaseous benzene, for example will fluoresce upon excitation in the UV  $(\lambda_{ex} > 260 \text{ nm})$  at a wavelength around 300 nm. If benzene is frozen in a glassy matrix, phosphorescence is observable around 340 nm.

Furthermore, the molecule can undergo an irreversible change, causing it to cease fluorescing ( $k_b$ , "bleaching" or "photo-destruction"). Concerning the exact nature of this photo-destruction, literature is not completely clear. Some attribute this to a process originating from the excited singlet state[22], but most sources consider bleaching to be a process involving the triplet state[23], coupling to the triplet state of oxygen, which will form a highly reactive singlet oxygen which may destroy the chromophore. This latter model is the one implied in figure 1.1.

<sup>&</sup>lt;sup>2</sup>For a more complete overview on these processes as well as other de-excitation pro-

	eYFP	TMR
$\sigma \cdot 10^{-16} \ \mathrm{cm}^2$	2.5	1.9
$k_s\cdot 10^9~{\rm s}^{-1}$	0.27	0.48
$k_{ISC}\cdot 10^6~{\rm s}^{-1}$	*	14.2
$k_t \cdot 10^6 \mathrm{~s}^{-1}$	*	0.5
$k_b \cdot 10^3 \mathrm{~s}^{-1}$	0.29	0.14

Table 1.2: Rate constants for some fluorophores. the values for eYFP are taken from [23], those for Tetramethyl Rhodamine from [17]. (\*) For eYFP the ratio  $f_t = k_T/k_{ISC} = 0.08$  is known.

$$\frac{\partial N}{\partial t} = \begin{pmatrix} -k_{exc} & k_S & k_T & 0\\ k_{exc} & -k_S - k_{ISC} & 0 & 0\\ 0 & k_{ISC} & -k_T - k_B & 0\\ 0 & 0 & k_B & 0 \end{pmatrix} \cdot N, N = \begin{pmatrix} N_G \\ N_E \\ N_T \\ N_B \end{pmatrix}$$
(1.3)

The kinetics of such a four-level system can be described by the differential equations given in 1.3. They can be solved numerically with small effort. Typical values for the rate constants are given in table 1.2. As the normalization condition  $\sum_{j=G,E,T,B} N_j(t) = 1$  applies, the system is of rank three. The general solution will be a superposition of exponential functions  $N_j = \sum_{i=1..3} C_{j,i} \cdot exp(-\lambda_i t)$ , where  $\lambda_i$  are the three eigenvalues of the matrix A and  $C_{j,i}$  are constants determined by the normalized eigenvectors and initial conditions. While literature mentions closed solutions to this



cesses, please refer to Valeur [21].

system[see e.g. 22, 24, 25], it is usually assumed that bleaching is small or negligible, resulting in an easily solved second order system. The complete third order system of equation 1.3, however, can be solved analytically only by applying Cardano's formula, which yields rather complex results which can not be simplified, making it very difficult to interpret them.

In the case of a small bleaching rate the fluorophore can be treated almost as a three-level system insofar that the population ratios of ground, excited and triplet state are close to the steady states of the simpler threelevel system.

For the autofluorescent proteins, however, this simple model has severe limits, caused by the complicated photophysics of e.g. GFP and its variants[26]. In addition to the states mentioned before, the autofluorescent proteins undergo changes associated with protonation of the chromophore, causing them to either shift in wavelength or go into one of a number of extended "dark states", remaining "silent" or unexcitable for times of a few microseconds to some hundreds of milliseconds. This also leads to the interesting observation that the mean bleaching time seems to be up to two orders of magnitude larger in bulk than for single-molecules[23], as the single molecule microscopy tends to lose track of molecules that stay dark for more than one image, typically less than 100 ms. Fluorescence Correlation Microscopy (FCS)[27, 28] has been used to show that this "blinking behavior" of autofluorescent proteins is present on different timescales.

# 1.3 Advantages and Shortfalls of Single Fluorophore Markers

#### 1.3.1 Shortfalls

One of the most important limitation imposed on measurements with single fluorophores as markers is photo-bleaching, limiting the number of photons which will be emitted during the lifetime of a fluorophore.

$$F_{sm} = \frac{k^{\infty}\tau}{1 + \frac{I_S}{I_{ill}}} \cdot \left(1 - e^{-t_{ill}/\tau}\right)$$
(1.4)

Equation 1.4 gives the number of photons emitted,  $F_{sm}$ , as a function of the illumination time  $t_{ill}$ , the bleaching time  $\tau$ , the maximal emission rate  $k^{\infty}$ , and the intensities of illumination  $(I_{ill})$  and saturation  $(I_S)$ . The bleaching factor  $(1 - \exp(-t_{ill}/\tau))$  can be approximated to  $t_{ill}$  if the experimental conditions are chosen such that  $t_{ill} \ll \tau$ . This leads to the well known linear relation between collected signal  $F_{sm}$  and the illumination time.

In the case of autofluorescent proteins, and for the limit of individual fluorophores, the mean time before photobleaching is almost equal to the typical time of illumination[23]. The effect of bleaching during illumination can therefore not be neglected.

As mentioned before, Harms et al. [23] also showed that the bleaching time of fluorescent proteins appears to be higher when measured in bulk  $(\tau_{bulk} = 112 \pm 3 \text{ ms})$  than when measured on the level of single molecules  $(\tau_{sm} = 3.5 \pm 0.5 \text{ ms})$ . A viable explanation for this seemingly inconsistency is the phenomenon of "blinking", a process in which the fluorophore



converts into a long-lived "dark state" when excited. For all practical applications, we are not able to distinguish between a photo-damaged (bleached) molecule and a long-term inactive ("blinking") fluorophore.

#### 1.3.2 Advantages

The use of single-molecule markers allows biological processes to be studied at early stages, when only a small number of molecules is involved. Many signaling cascades will eventually produce a reaction which can be observed with simpler optical methods (i.e. a morphological change of the cell) or by chemical means (i.e. secretion of a substance).

Most processes, however, have a discrete starting point, like the recognition of a hormone by a membrane-bound receptor or the opening of a membrane channel in reaction to an applied voltage. It is of great interest to study these processes in these early moments, as chapter 3 explains for a G-protein-coupled receptor.

Furthermore, the study of molecules one at a time makes a synchronization unnecessary. Interactions will be observed as they happen, regardless of when they happen. This also allows easy access to possible subpopulations, which can be detected even if they are small. Unlike bulk experiments, the signal from a small population is not overwhelmed by the majority population in the single-molecule approach.

Finally, many cellular proteins occur in small and tightly regulated numbers. Every deviation, both above and below the normal level, can cause severe repercussions for the biological system. Studying individual, labeled proteins allows us to maintain this biologically relevant expression level and avoids the massive over-expression needed for many biochemical assays.

## 1.4 Producing Fluorescent Samples

As the molecules of interest generally do not exhibit usable amount of fluorescence themselves, one has to label them before measurement[29]. There are many different ways to attach flourescent labels. In the following I will introduce several of the possibilities.

## 1.4.1 Artifical Fluorophores

Artificial fluorophores come in a large variety of colors and with various protocols that allow the controlled attachment of the fluorophores to biomolecules.

Figure 1.5 show two examples of such fluorophores, tetramethyl rhodamine (TMR) and 4,4-difluoro-4-bora-3a,4a-diaza-s-indacene (BODIPY), which stand as generic models. By small modifications of the sidegroups, these fluorophores, as well as other, commercial fluorophores like Alexadyes or Cy-dyes, can be adapted for different wavelength regimes over the whole visible spectrum. The attachment process can be either covalent, as shown in figure 1.5C&D, where the widely used labeling of free sulfur and primary amines are shown. It has to be mentioned that these reactions again display a multitude of variations (addition of hydrocarbon spacers, modifications of the side chains) which will influence the specificity of the labeling, preferentially targeting certain sites in biomolecules.

Besides the direct labeling of biomolecules, one can also utilize noncovalent binding with antibodies, immunoglobulins (Ig), the specific interactions of avidin and biotin (biotinylated fluorophores) or nitrilotriacetic acid (NTA), which will chelate with a motif formed by six histidine amino



acids  $(his_6-tag)$  in the presence of nickel.

The advantages of labeling with artificial fluorophores – availability of protocols for many different biomolecules, high signal and stability of the fluorophore – far outweigh their disadvantages for most applications. Yet especially in the limit of single-molecule fluorescence, there are significant problems with these labeling schemes: Most biomolecules present more than one possible attachment site, rendering the labeling a statistical process. The labeling of the protein can also lead to loss of function. Furthermore most protocols require an excess of the label to ensure a satisfying labeling yield; this excess has to be removed before measurement if individual labeled entities are to be identified. It is also necessary to test that the presence of an excess of unbound, free fluorophores does not lead to unspecific binding.

#### 1.4.2 Genetic Modification

As most proteins contain more than one possible labeling site, a method to increase the specificity of labeling is to use mutants of the proteins in which all but one labeling site are removed, or exactly one specific site is added. Obviously, such mutations have to be tested for their biological function. If several mutants are produced, which differ only in the position of the attached fluorophore, complex biological mechanisms, such as the function of motor proteins, can be studied[20] in great detail.

To obtain the protein, one can directly utilize the machinery inside a living cell. In a process called "transfection" a circular piece of DNA (plasmid) is introduced into the cell by means of a transfection agent. The plasmid contains the code of the mutant protein one wishes to study, a promoter sequence which causes the cell to transcribe the DNA, and a gene encoding resistance to certain antibiotics, so that one can selectively kill all cells which do not contain the plasmid. After a few hours, the cell will start to produce the protein. Most proteins are purified from bacteria or yeast cells, as those are easier to grow in culture than mammalian cells.

This can also be taken one step further by not only changing a few amino acids but creating a "fusion protein", a chimera of a small protein marker and a larger target protein. Again it is fundamental to ensure the functionality of this fusion protein. This method has become very important in recent years following the discovery of proteins which intrinsically exhibit fluorescence at a high level.

### 1.4.3 Autofluorescent Proteins

Already in 1962, Shimomura et al. [30] described the fluorescent properties of proteins found in the hydromedusa *Aequorea victoria* (figure 1.6). This finding was merely a footnote in their work to study the better known protein aeqorin, a chemoluminescent protein, found in this jellyfish. The emission spectrum of the wild-type GFP peaked at 508 nm.

The break-through for the autofluorescent proteins<sup>3</sup> came in with the cloning of the gene by Prasher et al. [32] in 1992 and the demonstration that GFP could be expressed in other organisms and still be fluorescent[33, 34].

Structurally, all fluorescent proteins derived from GFP share a motif of eleven  $\beta$ -sheets (forming the " $\beta$ -barrel") with  $\alpha$ -helices and coiled regions capping the structure (see figure 1.2). Another  $\alpha$ -helix lies inside this barrel. In the wild-type GFP, three amino acids in the center of this helix, <sup>65</sup>Ser-<sup>66</sup>Tyr-<sup>67</sup>Gly, form the chromophore *p*-hydroxybenzylideneimidazolin-

 $<sup>^3{\</sup>rm For}$  a comprehensive review on the Green Fluorescent Protein, see also the review by Tsien [31]. -13 –



#### CHAPTER 1. INTRODUCTION

one[33] after the protein obtains its tertiary structure[31]. This process is also referred to as "maturation" of the Fluorescent Protein and occurs within hours if oxygen is present. One consequence of the oxidative formation of the chromophore is the release of hydrogen peroxide, which could explain that high-level expression of GFP can affect the viability of cells.

To overcome limitations of the wild type protein, especially its tendency to dimerize, to yield low fluorescence, and to exhibit a strong dependence on the pH of the surrounding medium, "enhanced" or eGFPs were engineered by mutagenesis of the original wild type gene. Substitutions of the chromophoric amino acids as well as surrounding groups led to a variety of mutations with different absorption and emission properties. Today four groups of fluorescent proteins derived from the original *Aequorea*-GFP are available:



Figure 1.2: Molecular Model of GFP; the 11  $\beta$ -sheets forming the barrel structure are shown as bands. The chromophore is located inside this barrel.

blue (eBFP; excitation: 360 nm, emis- Source: Protein Data Bank, 1EMB sion: 442 nm), cyan (eCFP; excita-

tion: 452 nm, emission: 505 nm), green (eGFP; excitation: 488 nm, emission: 508 nm), and yellow (eYFP; excitation: 514 nm, emission: 527 nm). The actual number of distinct mutations is probably already close to a hundred.[data taken from: 31]

A recent addition is is the even more red-shifted DsRed, a GFP-like protein derived from the tropical coral *Discosoma sp* (figure 1.7), which absorbs around 560 nm. While the application of this red variety has been limited due to its inherent formation of tetramers, enhanced, monomeric versions (mDsRed) have been described.[35]

## 1.5 Scope of the Thesis

# 1.5.1 Application of Fluorescent Proteins with Two-Photon Excitation

In Chapter 2, the properties of the autofluorescent proteins are examined with two-photon excitation and compared to that of flavin. Flavin, while a weak fluorophore, is part of the cellular energy cycle and therefore highly abundant in living cells. In this study, flavin is used as a model for the cellular autofluorescence.[36]

Two-photon excitation is a non-linear process, in which a fluorophore absorbs two photons in rapid succession, each of which carries half the energy needed for the transition between ground state and excited state. The (informal) unit for the two-photon absorption cross-section is named after Maria Göppert-Mayer (1 GM =  $10^{-50}$ cm<sup>4</sup> s/photon), who predicted this process in 1931.[37]

Compared to the one-photon absorption cross-section ( $\sigma = 10^{-16} \text{ cm}^2$ ) the two-photon absorption cross-section is very small. Even as the probability for a two-photon event is dependent on the square of the laser intensity, one still needs an about six orders of magnitude higher intensity to obtain the same probability of absorption. Important is the num-



ber of photons of wavelength  $\lambda$  per Joule, which can be calculated by  $N_{p\lambda} = \frac{\lambda}{hc} = \lambda \cdot 5.03 \cdot 10^{15} \ (\lambda \text{ in nm}).$  A typical excitation intensity for one-photon fluorescence measurments is 1 kW/cm<sup>2</sup>, which corresponds to  $N_{p, 514\text{nm}} \simeq 2.6 \cdot 10^{18} \text{ photons/ms/cm}^2.$ 

These comparatively large laser intensities can be achived by a pulsed laser system which will generate short (100 fs) pulses with a very high peak power.

The basic advantage of two-photon excitation is based on the earlier observation that the two-photon absorption cross-section of a variety of fluorescent molecules scales super-linearly with the one-photon absorption cross-section[38]. Furthermore, due to the low absorption of infra-red radiation in biological materials, it is possible to image deep into tissue with negligible damage.

It was the hope to use autofluorescent proteins as single-molecule markers in *in vivo* two-photon-microscopy, which would have allowed to greatly reduce the background usually plaguing measurements. All autofluorescent proteins exhibited large two-photon action cross-sections  $\sigma_{TPE}$ , ranging from 8 GM (eCFP) to 40 GM (eGFP, see table 2.1 for full details), and unlike in the one-photon case all fluorescent proteins were clearly distinguishable from the much smaller signal of flavin ( $\sigma_{TPE} = 0.4$  GM). In conclusion, however, the rapid photo bleaching proved prohibitive[39] and does not allow their use as single-molecule markers under two photon excitation.

#### 1.5.2 Single-Molecule Studies of a Membrane Protein

In chapter 3, stimulation of  $\beta_2$ -adrenergic receptors ( $\beta_2$ ARs) is studied

on the single-receptor level. The  $\beta_2 AR$  is a member of the family of Gprotein-coupled receptors (GPCR, figure 1.3), which share a hepta-helical structure, are the largest class of cell-surface receptors, and are a major target for drug development.[for an overview see: 41] The  $\beta_2 AR$  is a receptor found in many different tissues and it mediates the action of catecholamines – neurotransmitters which are derived from tyrosine, such as dopamine and epinephrine – that produce relaxation of smooth muscles and vasodilatation.



Figure 1.3: A model of a G-protein coupled receptor (GPCR). Common to all members of this large family are the seven transmembrane helices. Figure adapted from [40]. Especially the role of oligomerization of the receptors is a topic of great interest[42], as the exact function and mechanism behind the formation of homo- and hetero-dimers between closely related GPCRs in live cells is still under debate.

To study the oligomerization of the receptor with our highly sensitive single-molecule methods, we produced a  $\beta_2$ AR-eYFP fusion protein in collaboration with the Department of Medicinal Chemistry at the Vrije

Universiteit Amsterdam. This construct was tested and found to have fully retained its biological function.

We set out to test the hypothesis that monomeric receptors are recruited into functional dimers upon stimulation. As the stoichiometry of individual entities on the cell membrane changes, their intensity will change



accordingly.[17]

To quantify the resulting intensity distributions (figure 3.2) we had to characterize the signal distribution from known monomeric YFP-fusion proteins and develop models to account for the bleaching and blinking behavior of the fluorescent proteins.

Computer simulations show that this is no trivial task, as for higher order aggregates (> 3) the intensity distributions of the oligomers start to smear out as shown in figure 3.4. It deviates markedly from the ideal of a multi-peaked distribution that can be found in model systems and for fluorophores which are stable longer than the average time of illumination.

Our measurements show that the  $\beta_2$ AR can be found in a mostly dimeric state before stimulation and that it will form larger aggregates within less than one minute after simulation. These aggregates remain stable on a timescale of minutes.

## 1.5.3 Technical Application of Single-Molecule Microscopy in DNA Research

In chapter 4 an application for rapid and separation-free detection of single base-pair mismatches of DNA is presented.

The method utilizes the high sensitivity of our setup to detect a strand of multiply labeled DNA against the background of excess labeled probe.

In contrast to existing methods like PCR (Polymerase Chain Reaction), which can also be used to detect variations in only one DNA base, this new technique does not require temperature cycling or cleaning steps to remove unbound probes.

This is made possible by Rolling Circle Amplification [43, 44] of spe-

cially designed "padlock"-probes, which generates a long, repeated, singlestranded piece of DNA upon recognition of a certain DNA sequence. The recognition step in this reaction is sensitive enough to be blocked by a single basepair mismatch. Figure 1.8 gives a schematic presentation of this technique.

The recognition site of the padlock probe is located on the ends of the DNA strand. Only if the template strand is matched perfectly a ligase will be able to ligate the two ends of the padlock, brought together by the recognition of the target. A circular piece of DNA is formed in this process. Addition of  $\Phi 29$  DNA polymerase will generate the long, repeated sequence of single-stranded DNA mentioned above.

With repeat-length of up to 1000 copies, binding of fluorescently labeled oligo-nucleotides resulted in an extremely high localized signal, which was detected against a 1000-fold excess of free fluorescent probes.

In addition to the adaptation of the optical setup to allow the observation of a larger area ( $100 \times 100 \ \mu m^2$ ), a flow system had to be implemented to screen larger amounts of sample in a fast, controlled and thereby reliable way.

## 1.5.4 Simultaneous Imaging and Spectroscopy

As the previous chapters concerned mainly the quantification and localization of single molecules, chapter 5 covers work on simultaneous localization and spectroscopic analysis of small, dilute fluorescent particles using a reflective blazed grating.

The technique described in this chapter was used in two distinct ways. It presents an enhancement of previously described methods[45], in which



#### CHAPTER 1. INTRODUCTION

two different fluorophores could be simultaneously observed using a wedge mirror. By using a blazed grating, I generalized this to an arbitrary number of fluorescent labels, only limited by the discrimination permitted from the resolution of the grating and the requirement of simultaneous excitation of the fluorophores.

One venue opened by this imaging spectroscopy method is the study of photosynthetic Light Harvesting Complexes at 4.2 K. Current practice is the identification of these LH-complexes by widefield imaging, while the subsequent spectral analysis has to be done with confocal methods[46]. By using a grating, one can combine these steps, simultaneously obtaining the emission-spectra of a number of complexes.

In our experiments we characterize the sensitivity of the method by using 20 nm fluorescent latex beads with spectra in the yellow ( $\lambda_{em} = 520$  nm) and red ( $\lambda_{em} = 600$  nm), each giving a signal equivalent to ~ 180 fluorescein molecules. The emission spectra gained from such individual beads are in ex-



Figure 1.4: The COBRA-FISH labeling for n = 3, r = 3 and one binary label. The number of possible combination is 24, allowing each chromosome in the mouse genome to be labeled with a unique combination.

cellent agreement with bulk-spectra obtained of beads in solutions (figure 5.2).

In a more application directed approach, we use this new technique to the karyotyping of mouse chromosomes. In a COBRA-FISH<sup>4</sup> [47] stained sample each chromosome carries a defined combination of four fluorophores in specific ratios, thereby allowing to uniquely identify each chromosome.

Conventionally, the chromosomes are identified by consecutive imaging with different excitation and emission filter set to identify the component fluorochromes. With our new technique, a single exposure is enough, to read out the full information.

<sup>&</sup>lt;sup>4</sup>**co**mbined **b**inary **ra**tio - **f**luorescence **i**n-**s**itu **h**ybridization. A technique which allows to define a combinatory labeling by using *n* fluorochromes, only two of which are used simultaneously per target, in *r* different intermediate labeling ratios, i.e. for r = 3 this would be 0% (none), 25%, 50%, 75% and 100% (full) staining. *m* additional "binary" labels can be added, resulting in  $n + \frac{r \cdot n!}{2(n-2)!} \cdot 2^m$  distinct color codes (figure 1.4).





Figure 1.5: Two examples and reaction schematics of commercially available fluorophores for the labeling of biomolecules. Tetramethyl rhodamine-6-maleimide (**A**) will form a thioether with free sulphur groups (**C**), such as can be found in cysteine-residues of proteins. BODIPY-FL (**B**) contains a succinimidyl ester that will form a carboxamide (**D**) with primary amino groups.(**E**) shows a non-covalent binding based on the specific recognition of the small molecule biotin (left) by the tetrameric protein avidin (right). Yet another common method to attach or immobilize is the use of NTA, which will form a stable chelate with a his<sub>6</sub> motif of proteins in the presence of nickel (**F**). Source for the model of Avidin: Protein Data Bank, SWE1.



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Figure 1.6: The hydromedusa *Aequorea victoria* naturally contains the original Green Fluorescent Protein.





Figure 1.7: The coral *Discosoma sp* possesses a protein "DsRed" which is very similar in structure to the green fluorescent protein from *Aequoria victoria*, yet exhibits red fluorescence.



Figure 1.8: The Rolling Circle Replication (RCR) depends on the recognition of the target DNA by a padlock probe. The recognition sequence is located at either end of the padlock, so that in the case of a match the probe can be circularized by a DNA ligase. A single DNA-mismatch will prohibit this ligation.  $\Phi 29$  DNA polymerase was used to form a long, single-stranded DNA consisting of repeats of the 93 base pair padlock probe.



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# Chapter 2

# Two-Photon Excitation Action Cross-Section of the Autofluorescent Proteins

The work described in this chapter has been published as "Two-Photon Excitation Action Cross Section of the Autofluorescent Proteins" in Chemical Physics Letters[1].

# 2.1 Abstract

We report on the values of the two-photon excitation action cross-sections of commercially available enhanced cyan, green, yellow and red fluorescent proteins. The two-photon absorption spectra are very similar in shape to those measured for one-photon absorption. However, they exhibit a significant blue shift, which is attributed to the participation of a vibrational



mode in the two-photon absorption process. The two-photon spectra are compared to that of flavine mononucleotide, which constitutes the main source of autofluorescence in mammalian cells. The definition of a relative detection yield between the autofluorescent proteins and flavine allows to quantify the applicability of autofluorescent proteins in two-photon singlemolecule studies in living cells.

# 2.2 Introduction

Major innovations in recent years have largely revolutionised the fluorescence imaging of biological samples. Two photon excitation using nearinfrared wavelength short laser-pulses increased the achievable penetration depth in thick samples, which for the first time renders imaging of thick tissues possible[2, 3]. In addition, the out-of-focus photodamage of the samples is essentially avoided which further facilitates reliable volume imaging of fluorescent samples.

Another advancement came from the discovery of the autofluorescent proteins originating from the jellyfish *Aequorea victoria*[4]<sup>1</sup> and the coral *Discosoma sp*[6]. By fusing a specific protein with an autofluorescent protein, the functional imaging of complex processes in cells became possible for the first time: the role of microtubuli in mitosis[7], functional oligomerisation of cellular proteins[8], or changes in the intracellular Ca<sup>2+</sup>-level[9, 10].

One advantage of the fluorescence labelling technique in comparison to bioconjugation is its extremely high sensitivity which makes experiments

<sup>&</sup>lt;sup>1</sup>For an overview of autofluorescent proteins see Tsien [5] and references therein.

feasible down to the level of a single molecule, opening a new avenue for cell biological and biophysical research. While the applicability of the autofluorescent proteins for single-molecule research *in vivo* has recently been demonstrated[8, 11, 12], its general application is still in early development. One major issue is the detectability of a single fluorophore within the autofluorescence background inside a living cell. The main component of this autofluorescent background are flavine molecules which are present at particularly high abundance of 100-1000 molecules per focal volume element[13]. Due to the broad absorption of flavine in the blue/green part of the visible spectrum the excitation spectra of the cellular autofluorescence strongly overlaps with the excitation spectra of the fluorescent proteins. For this reason, only the two most red-shifted varieties, the enhanced yellow fluorescent protein (eYFP) and the red fluorescent protein (DsRed), provide a significantly stronger wavelength discrimination compared to the flavines in wide-field single-molecule applications[11].

Here we report on the prospects of using two-photon excitation (TPE) in order to largely increase the signal-to-background ratio in our quest for general single-molecule detection of the autofluorescent proteins in live cells. The basic idea for two-photon excitation is resting on the earlier observation that the two-photon absorption cross-section of a variety of fluorescent molecules scales super-linearly with the one-photon absorption cross-section[14]. Hence, the ratio of the effective excitation rate of a fluorophore with high one-photon absorption cross-section, like the fluorescent proteins, and a fluorophore with a low one-photon absorption cross-section, like flavines, will be largely increased for two-photon excitation. Indeed our data show a super-linear scaling behaviour which has been previously de-



scribed only for small organic molecules. For exploitation and optimisation of the two-photon excitation scheme we have fully characterised and compared the two-photon spectroscopic properties of commercially available autofluorescent proteins and that of flavine mononucleotide. Our findings might lead to a novel strategy of two-photon imaging of single molecules using the autofluorescent proteins as markers.

# 2.3 Experimental Section

### 2.3.1 Sample Preparation

The autofluorescent proteins were purified as described previously[11]. In short, plasmids containing the coding sequences of the autofluorescent proteins with a C-terminal his<sub>6</sub>-tag (peXFP, Clontech) were transformed into E.coli and cultured at 37°C. The cells were harvested, lyzed and the fluorescent protein was extracted using a column of chelating sepharose (Pharmacia Biotech).

Concentrations of the fluorescent proteins were determined by measuring their absorption spectra. SDS-PAGE analysis revealed a correct molecular weight and an estimated purity of at least 95%.

25  $\mu$ l of stock solution of the fluorescent proteins and a stock solution of flavine mononucleotide were diluted in phosphate-buffered saline (PBS: 150mM NaCl, 15mM Na<sub>2</sub>PO<sub>4</sub>, pH7.4) to obtain a final concentration in the low  $\mu$ M range. Fluoresceine was diluted from stock into distilled water (NaOH to adjust for pH > 11).

All fluorophore concentrations were controlled spectroscopically.

### 2.3.2 Two Photon Excitation

Femtosecond pulses from a Ti:sapphire laser (Tsunami, Spectra Physics) with a mean power of 1-200 mW (150 fs pulsewidth, 80 MHz repetition rate) were focussed by a  $10 \times$  objective (Zeiss) into a quartz cuvette containing the solution. The focal spot had a diameter of approximately 100  $\mu$ m<sup>.</sup>

The fluorescence signal was discriminated from the excitation light by a 6 mm-thick blueglass filter (BG, Schott) and detected on an avalanche photodiode (SPCM-141, EG&G) under an angle of 90°. The overall detection efficiency of the entire setup was about 1%. Fluorescence signals were integrated over periods of 50 ms.

#### 2.3.3 Measurement of Two-Photon Action Cross-Sections

In this study we report on the two-photon excitation (TPE) action crosssection  $\sigma_{TPE}$ , defined as the product of the two-photon absorption crosssection,  $\chi^{(2)}$ , and the fluorescence quantum yield,  $\Phi$ .

All samples were measured at mean laser powers ranging from 1-200 mW (1-212 MW/cm<sup>2</sup> peak intensity) and up to 400 mW. The data were corrected for the background signal,  $F_0 < 10$  kcps, which was obtained from the pure buffer solutions on excitation.

Since the measurement of absorption cross-sections in the two-photon case heavily depends on the spatial and temporal coherence of the excitation beam[15] a calibration against a reference spectrum was employed using the routine described in detail inAlbota et al. [16]. In short, the fluorescence signal, F, was measured for the particular autofluorescent protein under study and compared to the fluorescence signal,  $F_{cal}$ , obtained from fluoresceine for which the two-photon action cross-section,  $\sigma_{TPE}^{cal}$ , has been reported[16]:

$$\sigma_{TPE} = \frac{\eta^{cal}}{\eta} \cdot \frac{c^{cal}}{c} \cdot \frac{F - F_0}{F^{cal} - F_0} \cdot \sigma_{TPE}^{cal}$$
(2.1)

The collection efficiencies,  $\eta$  ( $\eta$ =0.5% for DsRed,  $\eta$ = $\eta_{cal}$ =1.1% for all other fluorophores), and the respective sample concentrations, c, as obtained by one-photon absorption measurements on each sample were taken into consideration.

For wavelengths > 960 nm, where values for fluoresceine were lacking, the corrections were extrapolated from the 960 nm calibration value. The calibration procedure ensures that any effects such as the strong influence of spatial and temporal coherence of the excitation beam on the obtained results can be neglected.

## 2.4 Results and Discussion

The foremost characteristic of two-photon induced fluorescence is the dependence of the signal on the square of the excitation intensity[17]. For each wavelength measured we have recorded a power series of the detected fluorescence signal for peak intensities between 1–424 MW/cm<sup>2</sup> (figure 2.1).

For all 84 experiments we found that the fluorescence signal, F, followed a power-law dependence on the excitation power,  $F \propto I^{\alpha}$  with an average exponent of  $\alpha = 1.92 \pm 0.17$  (mean $\pm$ std). At the respective two-photon absorption maxima of the autofluorescent proteins, and thus at the highest signal-to-noise ratios, the average exponent was  $\alpha = 2.00 \pm 0.03$  (figure 2.1). Hence, it can be safely concluded that the detected signals originated from a two-photon excitation process. The TPE action cross-section spectra of the various autofluorescent proteins are given in figure 2.2 (data points) and are summarized in table 2.1. The maxima of the two-photon absorption spectra of the cyan (eCFP), green (eGFP), and yellow (eYFP) fluorescent proteins at  $\lambda_{max}$ =860, 920 and 960 nm, respectively, all fall into the tuning range of a standard Ti:sapphire laser equipped with broadband optics. The absolute values of the two-photon action cross-section,  $\sigma_{TPE}$ , at the maxima vary between 8 and 40 GM (1 GM = 10<sup>-50</sup> cm<sup>4</sup>s<sup>-1</sup>photon<sup>-1</sup>). The limited tuning range of our standard Ti:sapphire system did not allow for the identification of the maximum of the two-photon spectrum for DsRed, predicted to be around 1100 nm. The highest value obtained for DsRed was 20 GM at the wavelength of 980 nm.

It should be noted that the spectra and absolute values of the TPE action cross-sections reported here on eGFP are in good agreement with the values previously reported[3], and that of eYFP closely resembles the results obtained for a non-commercial, improved YFP mutant Citrine[18]. For the most redshifted autofluorescent protein, DsRed, current data suggest that it forms tetrameric complexes even at low concentrations[19] which, in turn, renders the definition of a molar absorption  $\varepsilon$  and thus the actual concentration in our experiment difficult[18, 20]. For the data presented the values of  $\varepsilon$  given by the manufacturer were used (table 2.1)[6, 21].

The increase of the TPE action cross-section at wavelengths below 760 nm, for which a sub-quadratic power dependence ( $\alpha = 1.68 \pm 0.12$ ) was found, does not correspond to any increase in the one-photon spectrum below 380nm. This new component has recently been described by Marchant et al. [22] as an irreversible "greening" of the protein by a three-photon





transformation process, which might also explain the deviation from the expected quadric power dependence. Such higher order photobleaching processes have been established before for a variety of other fluorophores[23].

Figure 2.2 is subsequently used for a quantitative comparison of the onephoton (solid lines, top wavelength axis) with the two-photon absorption cross-section spectra (points, bottom wavelength axis). The two-photon spectra roughly follow the shape of the one-photon spectra (when dividing the wavelength scale by two). However, a significant blue-shift in excitation is observed recurrently for the two-photon excitation spectra in comparison to the one-photon spectra[2, 16]. Whereas the absorption spectra are shifted, the fluorescence emission spectra stay unchanged for one and two photon excitation[24, 25], as predicted from Kasha's rule.

The blue-shift in the absorption spectrum is probably due to an additional, low-symmetry vibrational mode which couples to the electronic transition in order to get sufficient oscillator strength for the two-photon absorption process. For the pure electronic (0-0)-transition the two-photon absorption process is parity-forbidden. We quantified the blue shift for the autofluorescent proteins by a least-square fit which overlays the one- and the two-photon spectra. The values for the blue-shift found in this way are 255, 773 and  $1290 \text{ cm}^{-1}$  for eCFP, eGFP and eYFP, respectively.

It is interesting to note that vibrational modes of 222, 775 and 1507 cm<sup>-1</sup> have been identified for wild-type GFP by low-temperature line-narrowing experiments[26]. The different shifts found for the various autofluorescent proteins point to the appealing conclusion that depending on the fluorophore a different vibrational mode couples to the two-photon excitation process. However, it can not be excluded likewise that the different shifts

are mediated by a change in the protein environment.

The absolute TPE action cross-sections as determined for the autofluorescent proteins (figure 2.2) and the flavine mononucleotide (figure 2.3) determined here are used to quantify the relative detection yield of the autofluorescent proteins in comparison to flavines necessary for single-molecule studies in living mammalian cells.

Following a strategy developed by Harms et al. [11] we define the excitationwavelength dependent relative detection yield for the autofluorescence proteins as:

$$R_{(OPE/TPE)} = \frac{\eta}{\eta_{flavine}} \cdot \frac{\sigma(\lambda)}{\sigma_{flavine}(\lambda)}$$
(2.2)

where the experimental detection efficiencies for the various fluorophores,  $\eta$ , and the absorption cross-sections,  $\sigma_{OPE}$  and  $\sigma_{TPE}$ , for the one- or twophoton case are considered. In the case of two-photon excitation  $\eta/\eta_{flavine}$ was close to unity. The values for one- and two-photon excitation are summarized in table 2.2. In the one-photon case the low values of relative detection yield for eCFP and eGFP ( $R_{OPE} < 10$ ) render their use for wide-field single-molecule studies in cells extremely difficult. For eYFP and DsRed ( $R_{OPE} > 100$ ) single-molecule studies have been demonstrated[11].

This picture significantly changes for two-photon excitation. Indeed, due to the super-linear scaling of the two-photon absorption cross-section with the one-photon absorption cross-section (table 2.1), the relative detection yield for all autofluorescent proteins becomes  $R_{TPE} > 35$ . Such high relative detection yield leads to the extremely low autofluorescent background compulsory for single-molecules studies in life cells. Preliminary data from our laboratory show that the autofluorescent background is indeed largely (by a factor of ten or higher) reduced with respect to the autofluorescent background in one photon measurements.

# 2.5 Conclusions

We have shown that the use of two-photon excitation does largely enhance the detection ratio of all autofluorescent proteins against the autofluorescence background expected from flavines in living cells. The higher detectability is the result of the super-linear scaling between the one- and two-photon absorption cross-sections which has been previously reported for organic fluorophores. The presentation of the two-photon spectra allows one to optimize the excitation wavelength for any particular experiment including those in which one strives to use fluorescence resonant energy transfer between the various autofluorescent proteins[9, 10].

In conclusion, it should be mentioned that we have neglected the effect of photobleaching in our discussion, which is another factor important for single-molecule studies[11]. It has been generally found that the process of photobleaching is enhanced (typically by one order of magnitude) for two-photon excitation with respect to one-photon excitation[23, 24, 27] (see also figure 2.4). This further expends the enhancement of detectability due to the favorable spectroscopic properties investigated here. Further studies of the photobleaching behavior dependence on the pulse width and wavelength, for example, will be needed to finally decide if two-photon excitation will be superior to one-photon excitation for single-molecule studies in cells.

# 2.6 Acknowledgement

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fluorophore	one-photon excitation		two-photon excitation		
	$\lambda$ (nm)	$\sigma_{OPE}$	$\lambda$ (nm)	$\sigma_{TPE}$ (GM)	shift $(cm^{-1})$
		$(10^{-16} \text{ cm}^2)$			
eCFP	434	0.99	860	7.87	255
eGFP	489	2.10	920	41.21	773
eYFP	514	3.21	960	24.98	1290
DsRed	558	0.86	960	11.01	n.d.
FMN	266	1.22[30]	-		n.d.
	373	0.40	760	0.39	
	445	0.48	-		

Table 2.1: Maximal values of the one-photon absorption ( $\sigma_{OPE}$ , and the TPE action cross-sections ( $\sigma_{TPE}$ ; 1GM =  $10^{-50}$  cm<sup>4</sup> s/ photon) of the commercial autofluorescent proteins and of flavine mononucleotide (FMN). Blue-shift of the two-photon spectra with the respect to the one-photon spectra.



fluorophore	one photon		two photon	
	$\lambda ~({\rm nm})$	$R_{OPE}$	$\lambda \ ({\rm nm})$	$R_{TPE}$
eCFP	$435^{\dagger}$	3.4	860	36
	$457^{\ddagger}$	1.8		
eGFP	$488^{\dagger,\ddagger}$	8.7	920	467
eYFP	$514^{\dagger,\ddagger}$	405	960	566
dsRed	$532^{\ddagger}$	$> 10^4$	960	251

Table 2.2: Relative detection-ratio of the autofluorescent proteins versus flavine mononucleotide for one-  $(R_{OPE})$  and two-photon excitation  $(R_{TPE})$ . Values of  $R_{OPE}$  were taken from Harms et al. [11]. The value of  $R_{TPE}$  for DsRed will be considerably higher for its maximum absorption at wavelengths above 1000 nm.

(†: values for excitation at the peak absorption; ‡: values for wavelengths associated with the commonly used laser lines)



Figure 2.1: Dependence of the fluorescence signal on the excitation intensity for the autofluorescent proteins at their peak two-photon absorption wavelength; eCFP: 860 nm, eGFP: 920 nm, eYFP: 960 nm, DsRed: 980nm, and flavine mononucleotide: 760nm. The fluorescence follows a power-law dependence on the intensity with a mean exponent of  $2.00 \pm 0.03$  (dotted line).



Fig.1



Figure 2.2: Absolute TPE action cross-sections ( $\sigma_{TPE}$ , left axes) of eCFP, eGFP, eYFP, and DsRed in phosphate buffer (wavelength scale bottom). The values for  $\lambda \leq 960$  nm were calibrated against the values for fluoresceine (pH > 11)[16] (filled symbols). For  $\lambda > 960$  nm the calibration for 960nm was extrapolated. The one-photon absorption cross-sections ( $\sigma_{OPE}$ , right axes) are shown for comparison (solid lines; wavelength scale top).





Figure 2.3: Absolute TPE action cross-sections ( $\sigma_{TPE}$ , left axis) of flavine mononucleotides in phosphate buffer (symbols). The one-photon absorption cross-section ( $\sigma_{OPE}$ , right axis) is shown for comparison (solid line; wavelength scale top).



#### CHAPTER 2. TWO-PHOTON FLUORESCENCE



Figure 2.4: Two consequtive confocal images of tsa201 cells transfected with membrane-anchored eGFP (eGFP-C10HRas[28]). The images were taken with 5 ms illumination of 1.5 mW, 940 nm laser. The signal is comparable to the signal of this type of cell under one-photon illumination. About 90% of the fluorescence in the first image (A) is bleached after only one exposure (B). The bleaching is more than one order of magnitude faster than in the one-photon case[11].

As two-photon photobleaching is reported to depend on the third order [23] or even fourth order [29] of the excitation intensity, one would have to resort to long integration times ( $\gg 5$  ms) and low excitation intensity to be able to repeatedly image single molecules. This approach, however, does not allow for the study of dynamics inside the cell below the timescale of a few seconds.

Due to the limited intensity range used in preliminary measurements with spincoated Rhodamine we did not observe higher order bleaching as reported by other groups.[23, 29]

The scale bar is  $10\mu m$ .

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Chapter 3

# Aggregation of the $\beta_2$ Adrenergic Receptor Studied on the Level of Individual Receptors

The work described in this chapter will be published as "Single molecule studies of the  $\beta_2$  adrenergic receptor reveal aggregation before stimulation".[1]

# 3.1 Abstract

We present a study of the oligomerization of  $\beta_2$  adrenergic receptors ( $\beta_2$ ARs) upon stimulation with an agonist on live human cells. Using highly sensi-



tive microscopy techniques we follow individual eYFP-tagged  $\beta_2 AR$  ( $\beta_2 AR$ eYFP) molecules over time in the plasma membrane. Our data shows that the  $\beta_2 AR$ -eYFP is constitutively expressed as oligomers ( $N \leq 2$ ) and that further oligomerization (N > 3) of the receptor is induced within seconds after agonist stimulation.

## 3.2 Introduction

Guanidine nucleotide binding protein (G protein)-coupled receptors are encoded by the fourth largest gene family in man and play an important role in the transduction of extracellular signals across the plasma membrane of cells by the recognition and binding of specific ligands, such as ions, transmitters, and peptides, or – in the case of rhodopsin – by absorption of photons, upon which intracellular G proteins are activated. As such, these receptors mediate a variety of physiological functions, including vision, taste and olfaction, cognition, and immune regulation[2]. Until recently, GPCRs were thought to act as monomeric transmembrane entities. Accumulating evidence, however, now suggests GPCRs to exist in monomeric, dimeric, and various other oligomeric forms[3].

The  $\beta_2$  adrenergic receptor ( $\beta_2 AR$ ) is an important and well-characterized target for pharmaceutical drug development efforts as ligands for this receptor are successfully used for the treatment of a variety of conditions. The cloning of the  $\beta_2 AR$  gene showed that the  $\beta_2 AR$  and rhodopsin are members of a gene family that share sequence homology and a presumed "seven membrane spanning" topography [4]. A great number of receptors have been identified, and  $\beta_2$  adrenergic receptors have become the prototype for the intense study of GPCRs, which has lead to the unraveling of the structural basis of receptor function[5].

To date several techniques have been employed to study the dimerization of  $\beta_2$ ARs [6–8], including the immunoblotting upon immunoprecipitation [7, 9–11], the functional rescue by complementation of constitutively desensitized mutant receptors [12], computational methods [13], and a variety of fluorescent methods[9, 14–20]. GPCRs, including  $\beta_2$ ARs, have for instance been linked to fluorescent proteins with different fluorescence properties, such as the different fluorescent forms of the Green Fluorescent Protein (GFP) from *Aequorea victoria*, to allow the measurement of fluorescence resonance energy transfer (FRET) when two GPCRs are in close enough proximity [9, 14–17].

A widely used modification of this approach is bioluminescence resonance energy transfer (BRET) which has developed in recent years as a new technique to study protein-protein interactions. Protein partners of interest are tagged with either Renilla luciferase or green fluorescent protein (GFP). Non-radiative energy transfer between the excited luciferase and the GFP permits the study of spatial relationships between the two partners [9, 14, 18–20].

To investigate the potential oligomerization of the  $\beta_2$ ARs in real time we employed a highly sensitive microscopy method that allows the imaging and localization of individual eYFP-tagged  $\beta_2$ ARs in real time. Our studies show the  $\beta_2$ AR to be predominantly expressed as a dimeric receptor but the  $\beta_2$ AR is also constitutively present in the form of higher order multimers. Moreover, the  $\beta_2$ AR agonist isoproterenol modulates the relative ratios between the dimeric and multimeric forms of the  $\beta_2$ AR.



### 3.3 Experimental Section

### 3.3.1 Single Molecule Microscopy

The experimental setup has been described in detail previously [21, 22]. In short, HEK 293 cells were mounted onto an inverted microscope (Zeiss Axiovert100). The temperature was controlled by a stabilized water flow system and set to 37°C.

The sample was illuminated for 3 ms with the 514 nm line from an Ar<sup>+</sup>-laser (Spectra Physics), and imaged through a high aperture, 100× oilimmersion objective (Zeiss, NA = 1.4). The excitation light was suppressed in the detection path by use of a filter combination appropriate for eYFP (DCLP 530, HQ570/80, Chroma Technology; OG530, Schott). A 150 mm achromatic lens formed the image on a back-illuminated, liquid-nitrogencooled CCD camera (400 × 1340 pixels, 20 × 20  $\mu$ m<sup>2</sup> pixel size, LN/CCD-400-PB, Princeton Instruments). The total collection efficiency of the setup was better than 7%.

The camera was used in kinetic mode, resulting in sets of 10 images with a lag time of 5 ms between images. Each image was analyzed for single molecule signatures as described in detail previously[22]. In short threshold defined positions in the images were fitted with a Gaussian distribution, yielding positional information with 50 nm accuracy and the signal intensity with a precision of  $\approx 20\%$ , limited by shot-noise and the readout-noise of the CCD-camera.

### 3.3.2 Cell Preparation

HEK 293 cells were cultured in DMEM medium supplemented with streptomycin (100 $\mu$ g/ml), penicillin (100 U/ml), and 10% bovine serum in a humidified atmosphere (95%) at 5% CO<sub>2</sub> and 37°C. The cells were transferred every 4 days. To obtain expression of the  $\beta_2$ AR-eYFP fusion protein at the desired level, cells exhibiting confluence of approximately 50% were transfected with 1  $\mu$ g of cDNA and 3  $\mu$ l FUGENE-6 (Roche Biochemicals) six days prior to measurement. The cDNA was a gift from Dr. G. Milligan (Glasgow, UK), the construction of the receptor plasmide is described in Ramsay et al. [20]. Cells were separated and reseeded onto sterile 25 mm #1 microscope glass slides two days before measurement. The overall transfection efficiency was in the range of 30-50%. For measurements, cells were gently washed thrice with phosphate-buffered saline (PBS, 150 mM NaCl, 10 mM Na<sub>2</sub>HPO<sub>4</sub>, pH 7.4) before imaging in this buffer.

### 3.3.3 Stimulation Assays

For stimulation assays, cells exhibiting a low level of fluorescence were selected by eye and imaged as described above. After pipetting 5  $\mu$ l of 1 mM stock solution of either the agonist (–)isoproterenol or its biologically inactive isoform (+)isoproterenol (both: Sigma-Aldrich chemicals) into the cell-bath, cells were imaged for another 10 minutes. Due to the addition of liquid some refocusing was necessary, adding a delay of 1 minute between the images of unstimulated and stimulated cells.



### 3.3.4 Computer Simulations

Simulated intensity distributions were generated using Matlab (MathWorks, Inc.). Oligomers were considered to consist of N independent fluorophores with a normally distributed intensity of 200 counts ( $\sigma = 40$ ), each with a mean bleaching time of 1.5 images (4.5 ms) and a recovery probability of 1%, consistent with values published by Harms et al. [23]. The intensities of the active fluorophores where then combined into one peak, was placed into an image with a noise level of 6 counts per pixel (rms) and fitted with the same parameters as the measured data to ensure that possible influence of the fitting procedure would not affect the results.

The simulation of an oligomer ended when all N fluorophores were bleached.

### 3.3.5 Stoichiometric Analysis

Probability density functions (PDF,  $\rho(I)$ ) of measured and simulated data were fit to a model distribution using a non-linear least-square algorithm. The model distribution is derived from a "single molecule footprint" obtained by measurements of a membrane anchored eYFP moiety (hRASeYFP) expressed in mammalian cells.

The model allows for a shift in the peak intensity (maximum of the PDF) and generates oligomer distributions by convolution,  $\rho_{i+1}(I) = \rho_i(I) \otimes$  $\rho_1(I) = \int dx \rho_i(x) \cdot \rho_1(I-x)$  ("monomer-dimer-fit", figure 3.3 and table 3.1). The full distribution for *N*-oligomers,  $\rho_{total}(I) = \sum_N \alpha_i \cdot \rho_i(I), \sum_N \alpha_i = 1$ , is fully described by the monomeric distribution and the relative populations  $\alpha_i$ .

# 3.4 Results and Discussion

Heterologous expression of genes by standard transfection techniques generally results in the over-expression of the protein of interest. While this is of advantage for various biochemical assays, as it provides more material to be analyzed, it may also increase the probability of artifacts caused by the unphysiologically high density of the protein.

In contrast, highly sensitive analytical methods, such as single-molecule fluorescence microscopy, require low levels of protein expression, i.e. a low density of fluorophores. Transient expression of the  $\beta_2$ AR-eYFP in HEK 293 cells typically yields a receptor expression level of > 10<sup>4</sup>  $\beta_2$ AR molecules/cell, as can be determined by the fluorescence intensity of transfected cells (figure 3.1B). This number is two orders of magnitude above the single-molecule detection-limit of  $\approx 1$  fluorophore/ $\mu$ m<sup>2</sup>. We have previously used photobleaching to reduce excessive fluorophores before the measurement of single molecules [24]. However, this approach was considered not useful for the detection of receptor oligomers.

It is characteristic for transient transfection that, after an initial rise, the receptor-expression in transfected cells will diminish over time. For transiently transfected HEK 293 cells expressing the  $\beta_2$ AR we found that 6 days after transfection of the cells the average density of fluorophores on the cell membrane had dropped to a level where individual fluorophores can be identified (figure 3.1C). This low density ( $\leq 0.3 \ \mu m^{-2} = 0.013$  per pixel) also assures that the probability to find two receptors in close proximity purely by coincidence is less than 1%.

In order to analyze the local stoichiometry of the  $\beta_2 AR$  we also had



to establish a "single molecule footprint" for monomeric eYFP in a mammalian cell line. We did this using the membrane anchored hRAS-eYFP[25] under the same expression conditions described above for the  $\beta_2$ AR(figure 3.2, squares). This footprint was subsequently used for analysis as described in section 3.3.5. The intensity distribution of monomeric eYFP molecules that we obtained in HEK 293 cells is essentially identical to one measured for eYFP in artificial membranes or in gels[23].

The intensity distribution of  $\beta_2$ AR-eYFP deviates markedly from the intensity distribution we obtained for monomeric eYFP, indicating that against our predictions at least a part of  $\beta_2$ AR-eYFP are present as oligomers in cells at rest (figure 3.2, circles).

A further, dramatic change in the intensity distribution of  $\beta_2$ AR-eYFP is observed upon stimulation of the receptors with the  $\beta_2$ AR agonist (–)isoproterenol (figure 3.2,  $\triangle$ ). Analysis of the data indicates, in agreement with previous reports[14], that the  $\beta_2$ AR agonist promotes the aggregation of  $\beta_2$ AR-eYFP.

A change in stoichiometry can be quantified by a fit of the PDF according to a "monomer-dimer" model, in which one computes the expected di-, tri-, N-mer distributions by repeated convolution of the monomer distribution.

A more careful analysis of the data, however, shows that the intensity distribution obtained for the  $\beta_2$ AR can not be fully explained by the simple assumption that the previously mainly monomeric receptor now forms dimers. On the one hand, even before stimulation the intensity distribution deviates from the "single-molecule footprint". On the other hadn, a fit according to the "monomer-dimer" model described previously reveals that the intensity distributions are broader and show less pronounced features than expected (figure 3.3), especially in the case of the stimulated  $\beta_2$ AR.

The cause for this deviation is the particular bleaching and blinking properties of eYFP[26, 27], which we used as fluorophore for our experiments.

As the case of the negative control with (+)isoproterenol clearly shows, every population of oligomers will eventually bleach, leading to a strong monomer-like peak. In the case of both (+) and (–)isoproterenol stimulation, the cells were imaged before and after addition of the substance. Therefore the stimulated  $\beta_2$ AR-eYFP are expected to exhibit stronger effects of photobleaching.

This also means that the fitting method will severely underestimate the actual number and size of oligomers present. Photobleaching continuously removes fluorophores during the measurement, and in the case of stimulated cells, even before the measurement.

The number of fluorophores bleached is calculated by using the probability  $P(t) = \int_0^T \frac{1}{\tau} e^{-\frac{t}{\tau}} dt$  that a fluorophore with bleaching constant  $\tau$  will be bleached after time t. For N independent fluorophores in an aggregate, the chance that n of those are bleached is given by the binomial distribution

$$B_{p}^{N}(n) = {\binom{N}{n}} p^{n} (1-p)^{N-n}$$
(3.1)

where  $p = P(t_{exposure})$  denotes the probability for each fluorophore to bleach during the exposure. After the i<sup>th</sup> exposure, the distribution is given by  $B_{p_i}^N(n)$ , with  $p_i = 1 - (1-p)^i$ .

As the bleaching is photoinduced, it will occur during acquisition of the image, thereby further broadening the distribution. To complicate matters



further, eYFP is known to exhibit blinking on the millisecond timescale, in which case one can not only identify stepwise reduction of signal of an aggregate of receptors, but with a small probability also the stepwise return of the signal.

To assess the influence of these photophysical characteristics of eYFP, we chose to simulate the behavior of N-mers with the same properties as found in eYFP[24] and compute the resulting PDFs. The curves in figure 3.4 have to be considered as an extreme case, in which an N-mer is followed until all fluorophores are bleached. Under this assumption it is generally not possible to clearly resolve the N peaks in the distribution which correspond to the consecutive bleaching of the constituent fluorophores.

The fitting of these simulation results, presented in table 3.1, shows that rapid photobleaching – while causing a severe under-estimation of the oligomerisation state – still permits to follow and distinguish changes in intensity caused by aggregation.

The qualitative shift in the intensity distribution observed after stimulation with (-)isoproterenol (figure 3.2) was quantified using the simulation results. The ratio of monomer/dimer of 81%/19% observed before stimulation are close to those found for dimers in the simulation, leading to the assumption that the  $\beta_2$ AR forms mainly dimers if expressed in a mammalian cell system. After stimulation with its agonist, and despite photo-bleaching, the distribution shifts markedly towards higher intensities (78%/14%/8% for mono-, di- and trimers, respectively), indicating the formation of aggregates with  $N \geq 4$ . The bleaching of such an aggregate with N = 4 is shown in figure 3.5.

In the control, a stimulation with the biologically inactive stereoisomer

(+)isoproterenol, the intensity distribution returns almost to the monomer distribution (7% dimer). As all stimulation experiments are conducted after taking images of the unstimulated cells, this shift was attributed to photo-bleaching as opposed to a reaction of the receptor.

# 3.5 Conclusions

We have shown that  $\beta_2 ARs$  are constitutively expressed as oligometric structures, mainly monomers and dimers when transiently expressed at low level in human cells. Stimulation of  $\beta_2 AR$  expressing cells with the  $\beta_2 AR$  agonist isoprotection increases the size of these  $\beta_2 AR$  oligometric from dimers to aggregates with  $N \ge 4$ .

This change in aggregate size is a clear reaction of the receptor to stimulation with an agonist. While we cannot rule out completely that the aggregation is caused – at least in part – by the downregulation of the receptor, our data does not show evidence for internalization. In contrast, the biologically inactive stereoisomer of isoproterenol does not modulate  $\beta_2$ AR oligomerization, as only bleaching of the eYFP fluorophore is observed, indicating that stimulation and aggregation are functionally linked.

In consequence we were also able to show that eYFP can be used as a marker to study aggregation of membrane proteins *in vivo*, despite its rapid photobleaching, which can interfere with the interpretation of the data if not properly taken into account. Furthermore we find a tentative link between the dimerization of GPCRs and their biological function in the surprising discovery that a biologically inactive hH1-GPCR occurs in monomeric form if expressed in HEK 293 cells..

![](_page_69_Picture_6.jpeg)

# 3.6 Acknowledgment

The  $\beta_2$ -receptor-eYFP fusion protein was a gift from Dr. G. Milligan (Glasgow, UK).

		monomer-dimer fit
$\beta_2$ AR unstimulated	monomer	81%
	dimer	19%
	trimer	-
$\beta_2 AR$ stimulated	monomer	78%
(-)isoproterenol	dimer	14%
	trimer	8%
$\beta_2 AR$ stimulated	monomer	93%
(+)isoproterenol	dimer	7%
simulation dimer	monomer	91%
	dimer	9%
simulation trimer	monomer	43%
	dimer	57%
	trimer	< 1%
simulation tetramer	monomer	26%
	dimer	38%
	trimer	34%
	tetramer	< 1%

Table 3.1: Results from fitting the intensity distributions of  $\beta_2$ AR before and after stimulation. While in any case addition of the agonist causes a shift from the monomer population into higher aggregation states. Simulations have been performed to see how the photobleaching affects this distribution. See section 3.3.5 on page 56 for more information on the model used for fitting.

![](_page_71_Picture_3.jpeg)
#### CHAPTER 3. SINGLE MOLECULE STUDIES OF $\beta_2 AR$



Figure 3.1: HEK 293 cells were grown in culture (**A**, white-light image); 24 hours after transfection the cells exhibited strong fluorescence (**B**,  $> 10^4$  fluorescent molecules per cell). The transfection efficiency was roughly 30%. The number of fluorescent molecules decreases over time to a level where single fluorophores can be separated (**C**, 6 days after transfection). Using a fitting procedure as described in the text, one can obtain positions and intensities of individual molecules (**D**). The scalebars are  $5\mu$ m.



Figure 3.2: Intensity distribution of individual  $\beta_2$ AR aggregates in HEK 293 cells. The "single molecule footprint" for monomeric eYFP has been obtained by an inactive GPCR-eYFP construct imaged under the same conditions. Even before stimulation, the intensities recorded for  $\beta_2$ AR show a clear deviation from the monomer distribution. Upon stimulation with the agonist (–)isoproterenol a marked shift towards higher intensities and therefore higher aggregation is observed. Addition of the inactive isoform (+)isoproterenol exhibits no such effect. The lower intensities found in this case are most likely an effect of photobleaching.





Figure 3.3: Example of the fit of individual  $\beta_2 AR$  in HEK 293 cells before stimulation. The experimental distribution is wider than the single eYFP "footprint" found for monomeric eYFP in mammalian cells. This widening can be explained by the presence of dimers and the effect of photobleaching.



Figure 3.4: Intensity distribution simulation of purely mono-, di-, tri-, tetra-, and pentameric fluorescence complexes. The mean bleaching and recovery times were chosen to match those of eYFP[24]. See section 3.3.4 on page 56 for details on the simulation.





Figure 3.5: A multistep photobleaching events for  $\beta_2$ AR after stimulation with (–)isoproterenol (N = 4, symbols) is described well by the intensity levels found in a fitting of the PDF. The levels (lines) of intensity are not evenly spaced due to the asymmetry of the monomer intensity distribution. The event occured at the end of a kinetic-mode-cycle, therefore the drop to monomeric level and eventual complete bleaching are not observed.

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## Chapter 4

# Novel Single Nucleotide Pairing Assay using RCA

The work described in this chapter is published as "Homogenous detection of single rolling-circle replication products" in Analytical Chemistry.[1]

## 4.1 Abstract

We describe a simple and straightforward approach for homogenous and isothermal detection of individual rolling-circle replication (RCR) products, which represent individual padlock probe circularization events. The RCR products constist of tens of kilobases long, single-stranded tandem repeated copies of the probe sequence, and in solution they fold into micron-sized random coils. The method is based on the local enrichment of fluorescence labeled probes that hybridize to the coiled RCR products compared to the



concentration of free probes in solution. We present a detailed characterization of the fluorescence labeled products using a highly sensitive and fast microscopy set-up. At a  $10^4$ -fold excess of free label we were able to detect and follow individual RCR products at a signal-to-background-noise ratio of 27. This high signal to background ratio leaves room for analysis in a simple detection device, at higher speeds or at lower labeling ratios.

## 4.2 Introduction

Probe circularization reactions with padlock probes allow sensitive and specific detection of DNA and RNA sequence variants[2–6]. In these studies padlock probes have been visualized in situ with fluorescence labeled antibodies or they have been amplified with PCR for subsequent analysis of the amplification product on gels or DNA micro-arrays. Rolling circle replication (RCR) specifically amplifies circularized probes, but the method suffers from limited sensitivity due to the linear mode of amplification[7, 8]. RCR generates a product that is different from most DNA molecules present in a typical biological sample, since the RCR product is a single-stranded tandem repeated copy of the circularized probe, typically thousand units long, and readily available for hybridization. Analysis of individual RCR products, representing single molecule detection events, has been described for analysis of RCR products in a solid phase based assay[9]. The detection efficiency was limited, however, probably due to slow hybridization kinetics on the flat glass surface used in the study.

In this paper we present a method for homogeneous and highly sensitive detection of RCR products in the presence of a ten thousand-fold excess of unbound probes. The detection is performed in a microscope equipped with a flow system. The properties of labeled RCR products in solution are characterized, as well as the signal-contrast obtained between unbound fluorescence labeled probes and the ones bound to the RCR products. The figure-of-merit and the limitations of the novel method are discussed.

## 4.3 Experimental section

## 4.3.1 RCR template and probes

The RCR products were produced according to Nilsson et al. [8] using 180 fmol of the padlock probe pp93 (P-CCTCCCATCATATTAAAGGCTT-TCTCTATGTTAAGTGACCTACG ACCTCAATGCTGCTGCTG-TACTACTCTTCCTAAGGCATTCTGCAAACAT, P = 5' phosphate), circularized on 540 fmol of the ligation target GCCTTTAATATGATGGGA-GGATGTTTGCAGAATGCCTTAG, as polymerization substrate in a 90  $\mu$ l reaction volume. The  $\Phi$ 29 DNA polymerase-catalyzed polymerization reaction (90 min, 37°C) was terminated by 5 min incubation at 65°C. The RCR products were diluted to 1 pM concentration in PBS and labeled using different combinations of 10 nM 5' rhodamine- or Cy5- labeled oligonucleotide RC1 (CTCTATGTTAAGTGACCTACG), and/or 5' rhodamine labeled oligonucleotide RC2 (GCTGCTGTACTACTCTTCCT).

## 4.3.2 High sensitivity detection

To detect the amplified rolling circle products, we utilized a high sensitivity imaging setup, described in detail previously[10–12]. A 10  $\mu$ m thick



quartz precision cell (QS136, Hellma Benelux, Rijswijk, The Netherlands) was mounted onto a clean glass slide and subsequently fixed onto the microscope translation stage. The quartz cell allowed us to adjust a controlled flow through the detection volume and, in addition, confined the sample volume. Before adding the sample, the glass was coated with herring-sperm DNA to minimize unspecific binding of RCR products to the surface. While unspecific binding of the RCR products was effectively avoided in this way, a time-independent inhomogeneous background due to binding of detection oligonucleotides to herring-sperm DNA was observed. Because this background was static it was subtracted from the images, making identification of RCR-products straightforward (figures 4.4 and 4.5).

The microscope (Axiovert 100TV; Zeiss, Oberkochen, Germany) was equipped with a 10× Plan-Neofluar objective (NA = 0.3, Zeiss, Oberkochen, Germany). Samples were illuminated for 5 ms by the 532 nm line from a frequency-doubled Nd:YAG laser (Millenia X, Spectra Physics, Mountain View, CA, USA) for one color experiments. For two color experiments, the 514 nm line from an Argon-ion laser and 635 nm light from a tunable dye laser (model 3758, Spectra Physics, Mountain View, CA, USA) were used. The lasers were fiber-coupled into the microscope. A lens (f = 5 cm) focused the beam into the back focal plane of the objective leading to a lateral excitation profile of Gaussian shape with a width of  $w = 120 \ \mu m$ (full width at half-maximum). The mean illumination intensity, as calculated for the circular area of diameter w, was set to 0.2 kW/cm<sup>2</sup>. The illumination was accurately timed via a combination of mechanical shutter (Uniblitz, Rochester, NY, USA) and acousto-optic modulator (AOM, Isomet, Springfield, VA, USA). The use of appropriate filter combinations (one color experiments: dichroic 550DCLP and detection HQ 600/75, both Chroma Technology, Brattleboro, USA, in combination with OG570-3, Schott, Mainz, Germany; twocolor experiments: dichroic: TMR/Cy5 dual filter set, Chroma Technology and OG 570-3 Schott) permitted highly sensitive and specific detection of the fluorescence by a liquid-nitrogen-cooled slow-scan CCD camera system ( $20 \times 20 \ \mu\text{m}^2$  pixel size, Princeton Instruments, Trenton, NY, USA). A wedge-mirror simultaneously imaged two colors onto the surface of one CCD camera chip. The total detection efficiency of the experimental setup was about 1%. Images of  $50 \times 50$ pixel ( $100 \ \mu\text{m} \times 100 \ \mu\text{m} \times 7 \ \mu\text{m}$  depth of field) were taken every 100 ms, 500 images in sequence. An average background was computed from all images in a sequence and subtracted from each image before further analysis. To obtain accurate positional information and information on the signal level of the RCR products, we used a methodology developed in our lab for single-molecule tracking[13].

### 4.4 Results

The expected size estimated from the radius of gyration of a random coiled RCR product with 1000 copies of a 93 nt repeat is  $2 \cdot \sqrt{1000/6} \cdot 30$  nm = 775 nm, which is smaller than the pixel size of 2  $\mu$ m. Hence, the signal expected from an RCR product that has bound 1000 fluorophores (one/repeat), is located in one pixel of the CCD chip. In comparison, at a concentration of 10 nM of unbound detection oligonucleotide, 170 fluorophores are located within the detection volume of one pixel of a CCD chip (voxel) of  $2 \times 2 \times 7 \ \mu$ m<sup>3</sup>. Hence the RCR products are predicted to be



#### detectable at very high S/N ratio.

To test this prediction, padlock probes were circularized using an excess of target oligonucleotide ligation templates. The ligation templates further acted as primers in the subsequent 90 min RCR, which generates products containing 1000-1500 copies of the probe sequence[7]. The RCR products were diluted to 1 pM and subsequently incubated with 10 nM of fluorescence labeled detection oligonucleotides, complementary to a sequence present in each repeat of the tandem repeated single-stranded RCR products. Lower concentrations of detection oligonucleotide did not affect the background fluorescence significantly, and higher concentrations did not increase the signal significantly (not shown).

After subtraction of the static background signal (figure 4.5), individual RCR products appearing as signals with a lateral dimension  $< 2 \ \mu$ m, matching the point-spread function of the optical detection system, were clearly detectable even in the presence of a 10<sup>4</sup>-fold excess of free labeled oligonucleotides (figure 4.2A-C). The mean signal to background ratio seen in figure 4.2A-C amounts to S/B = 1 somewhat lower than the predicted value of 1000/170. However, relevant for identification is the mean signal to background-noise ratio, ( $S/N = S/\sqrt{B} = 27$ ) which was much larger than unity, rendering single RCR product imaging straight-forward. Negative control experiments were performed by omitting the incubation step with ligase. In those experiments we were unable to find any signals beyond noise after subtraction of the constant background signal.

Since RCR products appeared as point-like objects analytical methods developed for single-molecule studies were employed for further characterization. Individual RCR products were followed as they moved along with the flow of buffer through the detection volume (figure 4.2D). By their linear mobility they were clearly distinguishable from the inhomogeneous background. Further analysis of the trajectories revealed a flow speed of  $\simeq 10 \ \mu m/s$ , which translates to a screening volume of 0.5 nl/min. This speed of flow allowed accurate identification of the DNA particles in the solution with a camera readout speed of 10 images/s and illumination time of 5 ms. It should be noted, however, that the fairly low flow speed is not inherent to the method presented in this paper and can be increased to at least 100  $\mu m/s$  without smearing of signals within the 5 ms illumination time.

The signal level of the RCR products was further analyzed to obtain information about labeling efficiency and possible aggregation (Fig. 4.3). As the Gaussian illumination profile had a FWHM larger than the imaged area, no correction for changes in intensity was performed. We obtained mean signal values of  $637 \pm 64$  cnts/5 ms and  $858 \pm 86$  cnts/5 ms for RCR products labeled with rhodamine and Cy5, respectively. As the Gaussian illumination profile had a width larger than the imaged area, a correction for variation in intensity with position was omitted (figure 4.4).

Taking into account the expected signal from an individual fluorophore[10] of  $1.8 \pm 0.2$  cnts/5 ms and  $3.2 \pm 0.3$  cnts/5 ms for rhodamine and Cy5, respectively, at the present illumination intensity and detection efficiency, the RCR products generate signals that correspond to 300-350 pure fluorophores per RCR product. This number deviates substantially from the expected number of fluorophore-labeled oligonucleotides hybridizing to the 1000-1500 repetitions of the 93 nt motif. This discrepancy can be either due to incomplete hybridization, or more likely due to quenching caused



by electronic interactions between the fluorophores and the nucleobases, or with themselves at high packing. It has been shown from numerous studies that both rhodamine and Cy5 are sensitive to their local environment leading to self-quenching[14], base-specific quenching on binding to DNA[15] or even anomalous enhancement upon binding to proteins[16]. We have further shown that fluorescence from labeled oligonucleotides may be quenched by a factor of two when hybridized to intact RCR products, as compared to the same product in monomer form[8].

The intensity distributions for both detection oligonucleotides show features consistent with aggregation of two RCR products (figure 4.3). For the rhodamine labeled products a clear shoulder at 1301 cnts/5 ms is visible in the signal distribution (figure 4.3 left). The intensity distributions were further analyzed for two populations of single and double product formation utilizing a strategy developed earlier in our laboratory [17]. This method allows us to attribute an average stoichiometry to a whole population of products, even when the assignment my be difficult due to the large spread in intensities. Applying this strategy we found that  $47 \pm 5\%$ of the RCR products occurred as single products and  $53 \pm 5\%$  as double products for the rhodamine labeled products, and  $63 \pm 6\%$  "singels" and  $37\pm4\%$  "doubles" for the Cy5-labeled products. No separation-events of such "double" complexes were observed, leading to the conclusion that the complexes were stable over minutes. We have further noticed that aggregation of RCR products is promoted by high concentrations of BSA present in the polymerization buffer and can be avoided by modifying the buffer composition (J. Jarvius, personal communication).

## 4.5 Discussion

We have demonstrated the applicability of a homogeneous flow assay for the detection and analysis of 90 kb RCR products. Both rhodamine and Cy5-labeled RCR products show strong localized fluorescence, which was detected against the background of 10 nM labeled detection oligonucleotides, a concentration that seems sufficient to saturate the RCR products with labels.

Further, different fluorescence labels can be studied simultaneously, opening the possibility of parallel detection of several DNA sequences, or the further enhancement of detection sensitivity by coincidence analysis. In preliminary experiments we have applied a dual-color detection scheme[11] for detection of RCR products (figure 4.6). RCR products were labeled with both rhodamine and Cy5 labeled oligonucleotides at equimolar concentration. Clearly particles were identified in both channels with some of the particles showing fluorescence resonance energy transfer which would increase the specificity of this assay even further.

In general, analysis of individual single DNA sequences in complex DNA samples has been realized so far by detecting the coincident binding of two oligonucleotides labeled with two different fluorophores[18–20]. The present approach avoids problems of nonspecific background since detection events require coincident binding of many oligonucleotides to individual RCR products. Furthermore, the signal generated from hundreds of fluorophores allows faster detection with less advanced detection devices, compared to the set-up required for single fluorophore detection. With appropriate means of controlling the flow, as realized by microfluidics, the



screening of larger amounts (microliters) of solution within minutes will be possible, making our approach a viable means of biological and medical testing.

We expect that the method will present great advantages for multiplexed measurements of gene expression levels. Since differentially labeled RCR products are spatially resolved, a minority product will be detectable also in a large excess of other products, resulting in a high dynamic analysis range. Moreover, the quantitative precision should be superior to bulk measurements, since the precision of the present digital approach should ultimately be limited only by the statistical sampling error. We currently apply this methodology in a microfluidic set-up which can be mounted in either a conventional or a confocal fluorescence microscope (Jarvius et al., in preparation).

## 4.6 Acknowledgement

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Figure 4.1: Schematic representation of the setup. A shutter and an acousto-optic modulator (AOM) control the timing of the laser illumination; the 10  $\mu$ m high cuvette is homogenously illuminated and the emitted fluorescence is collected; a wedge mirror allows to split the image into a yellow (Rhodamine, top) and red (Cy5, bottom) channel, both of which can be imaged simultaneously. The cartoon on the left shows one 93 nt repeat of the DNA (blue) and the positions of the two probes (green and red, respectively). The DNA is not drawn to scale.



#### CHAPTER 4. HIGH SENSITIVITY DNA ASSAY



Figure 4.2: (A-C) Three images in a series of 54 images following a particular piece of DNA. The full trace is shown in black in all images. (D) Traces from 11 different pieces of DNA which could be followed for more than 30 images in one sequence (500 images); the trace from (A-C) is shown in red. The flow as calculated from the end-to-end distances of the traces is approximately  $10\mu$ m/s. (E) Mean background of the sequence, which has been subtracted before further processing. The scale bar in all images is  $10\mu$ m<sup>.</sup>



Figure 4.3: Distribution of fitted intensities for Rhodamine (left) and Cy5labelled DNA. A fit of both distributions indicates that 40-50% of the particles observed are actually two strands of DNA.





Figure 4.4: The fluorescent features caused by unspecific binding of oligonucleotide-probes to the herring-sperm coated surface stay constant in time (left). The mean background signal bleaches slowly over the course of 500 images (right top).

The analysis of intensities from individual RCR products (bottom right, dots) shows no strong correlation with the distance from the center of the laser profile (line: sliding average, triangles: standard deviation). The number of RCR products found at different distances from the center is given by the dark line.

The size of the image is 100  $\mu$ m; the laser profile is 120  $\mu$ m wide. Time per image is 5 ms.



Figure 4.5: The background subtraction uses the fact that the fluorescent features of the background are constant in time. The image and the movie show the raw data used to generate the traces in figure 4.2 (left), and the same image after background subtraction (right).

The size of each image is 100  $\mu {\rm m}.$ 





Figure 4.6: Comparison between RCR products labeled only with Rhodamine (left) and those labeled with Rhodamine and Cy5 (right). The excitation alternated between 635nm (Cy5 excitation; top) and 514nm (Rhodamine excitation; bottom). A 25% "bleed-though" of Rhodamines signal into the Cy5-channel is observed (bottom left and right). The signal from Cy5 was collected in the left channels, the signal from Rhodamine in the right channels. The images have been corrected for a slight vertical offset caused by the alignment of the wedge mirror.

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## Chapter 5

# Simultaneous Wide-Field Imaging and Spectroscopy of Localized Fluorophores

The work described in this chapter is published as "Simultaneous Wide-Field Imaging and Spectroscopy of Localized Flourophores" in Optical Letters.[1]

## 5.1 Abstract

A method that combines fluorescence imaging and spectroscopy of single molecules at room temperature is presented. This approach allows us to identify a number of imaged molecules unequivocally by simultaneously recording their fluorescence emission spectra. Furthermore, the spectral characteristics not only allow us to separate different fluorescent labels



quantitatively and qualitatively, but also provide information on the microenvironment of the molecules. This new method was successfully tested on a system of yellow-green and red fluorescent 20 nm latex beads and its benefit for the studies on biological systems was illustrated on a preparation of COBRA-stained mouse chromosomes.

## 5.2 Introduction

Advancements in optical methods and improvement of imaging equipment currently allow the routine application of ultra-sensitive fluorescence microscopy at the level of an individual molecule [2-4]. This development is particularly important in analytical chemistry and holds large promise for the biosciences. For the latter the detection and localization of molecular species present at very low concentration (a few copies per cell) is essential. Studies of molecular conformational dynamics, molecular interactions, and processes involved in signaling cascades are currently underway[5]. However, in multi-component systems like cells or biomolecular screening assays, one needs to go beyond mere detection and localization of one molecule to simultaneous, unequivocal identification of all components involved. The continued development of techniques to specifically label molecular species by selected fluorescence tags[4, 6, 7] allows the identification of a molecule by its fluorescence spectrum, which can also be used to obtain additional information about the local environment of the molecule. Hence, there is a general need for fast, parallel imaging technology with added spectral information. So far, technical solutions rely on confocal imaging[8], consecutive imaging with different emission filter sets[9], or on consecutive imaging with

different excitation wavelength[10, 11]. In a recent development, we have described a simple imaging scheme utilizing a dichroic wedge in the infinity path of a microscope to simultaneously acquire dual-color information on one CCD-camera[12].

A generalization of this scheme, which allows for simultaneous acquisition of image and spectrum, is presented here. In place of the dichroic wedge we mounted a blazed reflective grating into the infinity beampath of a microscope (figure 5.1). The dispersion of the grating was chosen such that both the zero'th and first order of diffraction were imaged onto the same CCD-camera. Thereby we obtained information on both the localization and the emission spectra of fluorescent objects which are sparsely distributed in the focal plane of the microscope. A similar approach has been described by Ma et al. [13], using a transmission grating which had the disadvantage that the intensity ratio between the zero'th and first order can not be adjusted. To test the applicability of our concept we used immobilized 20 nm fluorescent latex beads, and multiplex stained mouse chromosomes to demonstrate the biological applicability.

## 5.3 Experimental Setup

The experimental setup is based on the one described in detail previously[12]. In short, the samples were mounted onto an inverted microscope (Zeiss Ax-iovert100), imaged through a high aperture,  $100 \times$  oil-immersion objective (Zeiss, NA=1.4), and illuminated for 10 to 50 ms with a single line from an Ar<sup>+</sup>-laser, or at 625 nm from a dye laser (both Spectra Physics). The excitation light was suppressed in the detection path by use of an appropriate fil-



ter combination (chromatic beam-splitter, bandpass 514/560DBM, Chroma Technology; Holographic Super-Notch 488-1.0, Kaiser Optics; long-pass OG and RG filters, Schott). The blazed reflection grating (60 grooves/mm, 1.26° nominal blaze angle; Thermo RGL) was placed in the infinity path of the microscope (figure 5.1). A 150 mm achromatic lens focused the zero'th and first order reflections onto a back-illuminated, liquid-nitrogen-cooled CCD camera (400 × 1340 pixels, 20 × 20  $\mu$ m<sup>2</sup> pixel size, LN/CCD-400-PB, Princeton Instruments), forming an image and a first-order spectrum simultaneously.

### 5.4 Results and Discussion

#### 5.4.1 Calibration

A spectral calibration was carried out with the lines of the Ar<sup>+</sup>-laser, and the dye laser. The spectral separation d between the zero'th order and the first order for wavelength,  $\lambda$ , is determined by the grating equation,

$$d = f \cdot \tan\left(\arcsin\left(-\frac{\lambda}{g} - \sin\theta\right) + \theta\right) \tag{5.1}$$

where the focal length of the imaging lens, f, the angle of the incoming light,  $\theta \approx 45^{\circ}$ , and the grating constant, g. Equation 5.1 can be approximated as a linear relationship over the wavelength range of interest (500-700 nm) due to a small angular separation of  $< 4^{\circ}$ . The separation is characterized by the lateral dispersion of D = 1.5 nm/pxl. The blazing angle was chosen such that over the whole wavelength range of interest > 65% of the light was distributed to the first order. The total detection efficiency was better than 6% for all fluorophores used. The almost linear dispersion makes a calculation of the spectral image in the first order simple. Given the intensity distribution of the zero'th order image, I(x, y), being described by the sum of N separated molecular species  $I(x, y) = \sum_{i} f_i(x, y)$ , and the spectrum of each molecular species  $s_i(\lambda)$ , the first-order spectral image, S(x, y), is given by the convolution,

$$S(x,y) = \sum_{i}^{N} \int d\lambda f_i(x,y) \cdot s_i(\lambda)$$
(5.2)

There are two cases where analysis is most straightforward: (i) imaging of a single species (N = 1) of unknown spectrum for which  $s(\lambda)$  can be determined by image deconvolution, and (ii) imaging of multiple species of known spectra for which equation 5.2 can be used to identify signals in I(x, y) to originate from a certain species. Both approaches have been employed in our initial experiments.

#### 5.4.2 Nanometer-sized fluorescent beads

First, we tested our methodology on fluorescently labeled yellow-green and red 20 nm latex beads (Molecular Probes). Samples of both beads were spin-coated on HF-etched microscope cover slips to a surface concentration of  $< 0.1 \ \mu m^{-2}$  and were excited at 488 nm with 0.5 kW/cm<sup>2</sup> and 514 nm with 2 kW/cm<sup>2</sup>, for the yellow-green and red beads, respectively. Individual beads were detected at a signal-to-noise ratio of typically S/N = 30 in the zero'th-order image and the corresponding spectrum to each bead was identified in first-order (figure 5.2B).

Given the background noise of our camera system of b = 6 cnts/pxl (limited by the read-out noise), the width of a spectrum, w, of approximately



40 pixel (equivalent to  $\approx 60$  nm), and a total signal of S = 9000 counts collected for an individual bead, the signal-to-background noise ratio in a worst-case broad spectrum without features equals to  $S/B = S/(2 \cdot b \cdot w) =$ 19. In reality, however, the peak of a fluorescence distribution can exhibit an S/B ratio of up to one order of magnitude higher. At this high S/Bratio, deconvolution of the spectral image is straightforward. In the case of wavelength-delimited objects such as those shown in figure 5.2, the deconvolution step can be omitted completely. In either case this leads to the spectrum of the fluorescent beads given in figure 5.2B for the yellowgreen and red species, respectively. The resulting spectra show an excellent agreement with the corresponding bulk fluorescence emission spectra. This demonstrates the suitability of the new technique for obtaining fluorescence emission spectra of individual particles with very good spectral resolution and quality.

Each bead contains about 180 fluorophores (Molecular probes product sheet). Taking this and the above estimated S/B ratio into account, we extrapolate that we can obtain a quantifiable spectrum from the signal of just 10 molecules without further experimental refinement at an illumination time of just 50 ms.

#### 5.4.3 COBRA-FISH-stained mouse chromosomes

We also tested whether our methodology was able to localize and identify a species in a mixture of different fluorescent species. This is a typical problem encountered in chromosomal assays or karyotyping in which either colors are used to code for a specific chromosome (multicolor in situ hybridization (MFISH)[14, 15]), or in which a unique mixture of colors is used to code for the different chromosomes (combined binary ratio labeling in situ hybridization (COBRA-FISH)[16]). Here we studied COBRA-stained mouse chromosomes, carrying four different fluorescent labels (DEAC, FTIC, Cy3, and Cy5). Mouse fibroblast were arrested in metaphase by treatment with colcemide, swelled in hypotonic KCl solution, and fixed with methanol/acetic acid (3+1 v/v). This preparation was then dropped onto clean microscope slides, which resulted in random spreading of the chromosomes[17, 18] over an area of several  $\mu m^2$  (figure 5.3A).

In our experiment we concentrated on three of the dyes (FITC, Cy3,  $Cy_{5}$ ), by simultaneously exciting the sample at 514 and 625 nm. Figure 5.3 shows an image containing 9 separable chromosomes (figure 5.3A) and their spectra (figure 5.3B). In order to analyze the overlapping spectra in this experiment we utilized a non-linear fitting procedure to find the labeling ratios in each of the N = 3 dyes accessible for every chromosome. For this, chromosome locations were first identified by thresholding. Subsequently, equation 5.2 in combination with the normaized bulk spectra,  $s_i$  of FITC, Cy3, and Cy5 respectively, was used to generate a spectral image with the intensity values  $(I_{FITC}, I_{Cy3}, I_{Cy5})$  as free parameter for each chromosome (figure 5.3D). Fitting yielded the respective intensity values for each chromosome. The resulting chromosome identification image is shown in figure 5.3C. Not all chromosomes could be assigned uniquely, as we lacked information on the fourth label (DEAC) in the near UV. While the nonlinear fitting procedure works well for sparsely distributed species even in the case of overlapping spectral images (as shown in figure 5.3), this procedure will be increasingly error-prone for higher density samples. In this case, fluorescence labels with sharper spectra[19], leading to less overlap


in the spectral image, will be of advantage. Additional optimization and testing will be needed to assess whether our methodology will be applicable to high-density samples like DNA-chips.

## 5.5 Conclusions

In conclusion, we have presented a new method to simultaneously record fluorescence images and fluorescence emission spectra at high quality and spectral resolution, with an estimated sensitivity of about 10 single fluorophores. This new simple, robust, yet highly sensitive method can easily be incorporated in existing imaging setups. The speed of acquisition is limited only by the amount of signal that can be obtained from the sample and ranges from several images per second to video-rate and above. In turn, this lower time limit sets an upper limit to the speed of moving object which can be analyzed before smearing-out will occur in the image (at video-rate  $\approx 7\mu$ m/s). Indeed, any number of different fluorophores can be imaged simultaneously, provided that the emitting species are sparsely distributed over the field-of-view. We have shown that this methodology is suited for a variety of applications in life-science research.

### 5.6 Acknowledgments

Hans Tanke and Joop Wigand for providing us with a COBRA stained chromosomes. S. Oellerich acknowledges financial support from the European Marie Curie Fellowship program. Gregory Harms and Laurent Cognet are acknowledged for their valuable input at the early stages of this project.



Figure 5.1: (A) Experimental setup. The fluorescence image and spectrum of the sample are simultaneously formed on the CCD camera as the zero'th and first order diffraction from a blazed grating placed in the infinity beam path of the microscope and subsequently focused onto the chip. (B) Zoom into A; a wavelength calibration of the dispersed image is shown by the refraction of two out of six lines from an  $Ar^+$  laser; the inset on the right shows all six measured positions, resulting in a first order calibration.



Figure 5.2: (A) Image and Spectrum of yellow 20 nm beads. Starting from one bead one can accurately determine the distance between zero'th and first order, which are slightly offset in this example. The distance between zero'th and first order (arrow) can be used to determine a wavelength scale using the calibration shown in figure 5.1B. (B) Spectra obtained from individual yellow-green (left) and red (right) 20 nm beads (solid lines) in comparison with the corresponding bulk spectra (dotted lines).



Figure 5.3: Simultaneous illumination at 514 nm and 625 nm of metaphasic mouse chromosomes with COBRA-FISH staining. Several chromosomes can be identified in the zero'th  $(\mathbf{A}, 30 \times 30 \mu \text{m}^2)$  and first order image  $(\mathbf{B})$ . A least-square fit according to equation 5.2 (typical result in  $\mathbf{D}$ ) allows us to assign labeling ratios of the three different fluorophores to each chromosome (pie-charts in  $\mathbf{C}$ ). As the full labeling consists of four distinct colors, the chromosomes marked \* and # can not be uniquely identified.



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## **Summaries**

## Summary

### Single-Molecule Techniques in Biological and Biophysical Research

The rapid developments in the field of optical imaging and detection have led to a wealth of new methods which allow us to study matter on the molecular and even atomic level. While this makes it possible to visualize the structure of proteins and the organization of biological structures *in vitro*, and has provided answers to many issues, it has also raised new questions about the interaction of proteins with each other and their dynamics *in vivo*.

In my thesis I introduce several applications of single molecule techniques, mostly based on fluorescence microscopy. Chapter 2 concerns characterization of autofluorescent proteins in bulk for the case of two-photon excitation. Two-photon excitation is a non-linear process in which the energy needed to excite a fluorophore is taken from the absorption of two photons, each carrying half the needed energy, within a very short time ( $\approx 10^{-15}$  s). Despite the fact that this process is very rare, expressed in cross-sections,  $\sigma^{(2)} \sim 10^{-50}$  cm<sup>4</sup>s/photon versus the one-photon case



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 $\sigma^{(1)} \sim 10^{-16} \text{ cm}^2$ , it can be observed using a pulsed laser, generating ultrashort (100 fs) bursts of near-infrared light. The possible advantages of this method are the low absorption in tissue, allowing us to penetrate several millimeters into a specimen, and the low background expected due to the use of near-IR excitation wavelengths. Autofluorescent proteins proved to be excellent fluorophores in terms of two-photon absorption, yet their rapid photobleaching does not allow their use as single-fluorophore markers under these conditions.

Chapter 3 reports on true single molecule measurements using autofluorescent proteins as markers to study the aggregation of the  $\beta_2$  adrenergic receptor ( $\beta_2 AR$ ). The  $\beta_2 AR$ , a member of the large family of G-Protein-Coupled Receptors (GPCR), was one of the first GPCR to be sequenced and serves as widely accepted model system. It plays an important role in the regulation of blood-pressure and the dilation of the bronchia, and it can also be found in muscle tissue, the kidneys, the liver, and the pancreas. In biochemical assays GPCRs are regularly detected as dimers, and a connection between the formation of dimers and the functionality of the receptors has been suggested Our aim in this study was, therefore, to assess the degree of dimerization of the  $\beta_2$ AR-eYFP-fusion-protein in a mammalian cell prior to stimulation, and compare this data to the stoichiometry of the receptor after stimulation with the agonist (-)isoproterenol and its biologically inactive stereoisomer. To reach this goal we had to use the previous characterization of eYFP intensity and bleaching behavior to model the time dependent shift in intensity of oligometric receptors in membranes, and apply the result to the data found in our experiments. Indeed, our results show that the initially dimeric  $\beta_2 AR$  forms larger aggregates within one minute after stimulation. These aggregates stay stable for at least 10 minutes. The negative control with the biologically inactive stereoisomer caused a shift toward monomeric intensities, which can be explained by photobleaching.

Chapter 4 describes the adaptation of our highly sensitive single-molecule setup to allow the separation-free detection of Rolling Circle Replication (RCR) Products against the background of free, fluorescent labels.

RCR is a new and very specific way to test for the presence of DNA sequences. If a matching DNA-sequence is found by a so-called padlock-probe, this probe can form a circle. DNA-polymerase can copy this probe and, due to the circular form, will produce a long piece of DNA consisting of repeats of the padlock motif ("rolling"). Special labels recognizing this padlock motif can bind to the RCR product, causing a locally increased signal, which allows us to detect the RCR product without the need to first remove all unbound probes.

Finally, in chapter 5 a previous method allowing dual-color detection is extended to allow for simultaneous recording of images and emission spectra of small fluorescent particles. For this, a blazed reflection grating is introduced into the beam path. Unlike its conventional counterpart, a blazed grating allows one to choose and influence the distribution of energy between the different diffraction orders.

We applied this new technique to test samples of 20 nm fluorescent beads, membrane fragments containing photosynthetic complexes, and for the unequivocal identification of chromosomes using a COBRA-FISH<sup>1</sup> label.



<sup>&</sup>lt;sup>1</sup>COmbined Binary RAtio Fluorescence In-Situ Hybridization

### Samenvatting

### Technieken voor de observatie van individuele moleculen en hun toepassing in biologisch en biofysisch onderzoek

De voortdurende ontwikkelingen op het gebied van optische detectie hebben geresulteerd in een groot aantal nieuwe technieken, die wetenschappers toegang geven tot het niveau van enkele atomen of moleculen. Hierdoor is het mogelijk geworden de structuur van eiwitten en de organisatie van biologische systemen *in vitro* te visualiseren. Deze nieuw verworven kennis heeft veel vragen beantwoord, maar heeft ook nieuwe vragen over de dynamica en de interacties van eiwitten *in vivo* opgeroepen.

In mijn proefschrift introduceer ik toepassingen van technieken voor het visualiseren van individuele moleculen, die voornamelijk gebaseerd zijn op fluorescentie microscopie. Hoofdstuk 2 beschrijft de fotofysische karakterisering van autofluorescente eiwitten in bulk bij twee-fotonen excitatie. Twee-fotonen excitatie is een niet-lineair proces waarin de energie die nodig is om een fluorofoor aan te slaan afkomstig is van de absorptie van twee fotonen, waarbij beide fotonen de helft van de benodigde energie bijdragen. Dit proces vindt plaats op een tijdschaal van  $10^{-15}$  s. De waarschijnlijkheid van dit proces komt tot de uitdrukking in de absorptie doorsnede, die in het geval van twee-fotonen absorptie  $\sigma^{(2)} \sim 10^{-50}$  cm<sup>4</sup>s bedraagt, terwijl het in het een-fotonen geval  $\sigma^{(1)} \sim 10^{-16}$  cm<sup>2</sup> bedraagt. Het fenomeen van twee-fotonen excitatie kan geobserveerd worden door gebruik te maken van een laser die ultrakorte 100 femtoseconden pulsen uitzendt met een golflengte die in het nabij-infrarood ligt. De veronderstelde voordelen van deze methode zijn een lage absorptie in organisch weefsel waardoor het mogelijk wordt enkele millimeters in het weefsel door te dringen. Daarnaast is de verwachte achtergrond laag door het gebruik van excitatie golflengtes in het nabije infrarood. Autofluorescente eiwitten bleken fluoroforen van hoge kwaliteit te zijn waar het gaat om twee-fotonen absorptie, maar de snelle fotobleeking die optreedt maakt het onmogelijk ze te gebruiken voor studies op het niveau van enkele fluoroforen.

Hoofdstuk 3 behandelt een experiment toegepast op enkele moleculen. In dit experiment hebben wij autofluorescente eiwitten gebruikt als markers om de aggregatie van  $\beta_2$  adrenerge receptoren ( $\beta_2 AR$ ) te bestuderen. Deze receptor is onderdeel van de grote familie van G-eiwit-gekoppelde receptoren (GPCRs). De  $\beta_2$  adrenerge receptor was een van de eerste GPCRs waarvan de sequentie bekend was en dient als een algemeen geaccepteerd modelsysteem. De  $\beta_2 AR$  speelt een belangrijke rol in de regulering van de bloeddruk en verwijding van de bronchia. De receptor komt voor in verschillende weefsels zoals spierweefsel en organen zoals nier, lever en alvleesklier. In biochemische proeven worden G-eiwit-gekoppelde receptoren over het algemeen in de vorm van dimeren gevonden. Een verband tussen de vorming van dimeren en de functionaliteit van de receptor wordt al langere tijd gesuggereerd. Het doel van mijn onderzoek was dan ook de mate van dimerisatie van het  $\beta_2$ AR-eYFP fusie-eiwit in zoogdiercellen te bepalen, waarbij ook de invloed van stimulatie van de receptor door zijn ligand bekeken is. Om dit te bestuderen hebben wij gebruik gemaakt van de eerdere karakterisatie van de intensiteitsverdeling en het bleekgedrag van eYFP om de tijdsafhankelijke verschuiving van de intensiteit van geoligomeriseerde receptoren in membranen te modelleren. Het resultaat van deze modellering hebben wij gebruikt voor de analyse van de data die voort-



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kwam uit onze experimenten. Onze resultaten laten zien dat de, in eerste instantie als dimeer voorkomende,  $\beta_2 AR$  grote aggregaten vormt binnen een minuut na stimulatie. Deze aggregaten blijven tenminste 10 minuten stabiel. Proeven met een biologisch inactieve stereo-isomeer van het ligand als negatieve controle laten een verschuiving naar intensiteiten karakteristiek voor een monomeer zien. Deze verschuiving kan verklaard worden door het verschijnsel fotobleken.

Hoofdstuk 4 beschrijft de aanpassingen die wij aan onze opstelling hebben moeten aanbrengen, om de detectie van "Rolling Circle Replication" (RCR) producten mogelijk te maken tegen een achtergrond van fluorescente labels in oplossing, zonder zuivering van het monster te vereisen. RCR is een nieuwe en specifieke manier om de aanwezigheid van unieke DNAsequenties te testen. Als combinatie met een passende DNA-sequentie is opgetreden wordt het mogelijk voor de "padlock"-probe om enkelstrands, circulair DNA te vormen. Dit maakt het mogelijk voor DNA-polymerase om kopiën van deze probe te maken. De cirkelvorm bevordert de productie van lange stukken DNA ("rolling"), waarin de gezochte sequentie meerdere malen voorkomt. Speciale labels die deze sequentie herkennen kunnen aan het RCR product binden, waardoor lokaal een hoog signaal ontstaat, tegen de achtergrond van de ongebonden labels.

Ten slotte wordt in hoofdstuk 5 beschreven hoe de methode waarmee twee kleuren detectie kan worden toegepast, wordt uitgebreid om gelijktijdig zowel een plaatje als een emissie-spectrum van een klein fluorescent object te verkrijgen. Hiervoor wordt een *blazed* reflectie-tralie in het emissiepad geplaatst. Anders dan bij conventionele tralies is het met deze tralie mogelijk om de verdeling van energie tussen de verschillende diffractie-ordes te kiezen en beïnvloeden. Wij hebben deze nieuwe techniek toegepast op fluorescente bolletjes met een doorsnede van 20 nm, en membraan fragmenten die fotosynthetische complexen bevatten. Daarnaast kunnen wij met deze methode de verschillende COBRA-FISH<sup>2</sup> gelabelde chromosomen ondubbelzinnig identificeren.

## Zusammenfassung

## Einzelmolekül-Techniken in der biologischen und biophysikalischen Forschung

Die rasante Entwicklung auf dem Gebiet der optischen Mikroskopie resultierte in einer Vielzahl neuer Techniken, die die Betrachtung von Objekten auf der Größenskala von einzelnen Atomen und Molekülen erlauben. Dadurch ist es möglich geworden, die Struktur von Proteinen und die Organisation biologischer Systeme *in vitro* zu visualisieren. Die Kenntnis der biologischen Struktur half viele Fragen zu beantworten, hat aber zugleich weiterreichende Fragestellungen über z.B. die Dynamik von Proteinen *in vivo* und die Wechselwirkung zwischen Proteinen aufgeworfen.

In meiner Dissertation stelle ich einige Anwendungen dieser neuen Techniken zur Untersuchung von Einzelmolekülen vor. Die Techniken basieren hauptsächlich auf der Fluoreszenz-Mikroskopie. In Kapitel 2 beschreibe ich Experimente zur Charakterisierung von autofluroeszenten Proteine in Lösung unter Zwei-Photonen-Anregung. Die Zwei-Photonen-Anregung ist ein nicht-linearer Prozess, bei dem ein Fluorophor die Anregungsenergie E durch die Absorption von zwei Photonen mit der jeweils halben Ener-



 $<sup>^{2}\</sup>mathrm{COmbined}$ Binary RAtio Fluorescence In-Situ Hybridization

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gie E/2 erhält. Dies geschieht auf einer Zeitskala von etwa  $10^{-15}$  s. Die Wahrscheinlichkeit für diesen Prozess wird durch den "Absorptionsquerschnitt"  $\sigma$  beschrieben. Der Absorptionsquerschnitt für die Zwei-Photonen-Absorption beträgt  $\sigma^{(2)} \sim 10^{-50}$  cm<sup>4</sup>s, wogegen der Wert für "gewöhnliche" Ein-Photonen-Absorption bei  $\sigma^{(1)} \sim 10^{-16}$  cm<sup>2</sup> liegt. Durch den Gebrauch eines Lasers, welcher ultra-kurze (100 fs)<sup>3</sup>, intensive Pulse im nahen infraroten Spektralbereich produziert, ist es möglich, einen solchen Zwei-Photonen-Übergang zu induzieren. Die potentiellen Vorteile der Methode liegen in der geringen Absorption von infrarotem Licht, wodurch es möglich ist, einige Millimeter tief in intaktes Gewebe einzudringen. Des Weiteren ist bei dieser Anregungswellenlänge auch ein äußerst geringes Hintergrundsignal zu erwarten. Meine Studien zeigen, dass die autoflureszenten Proteine ausgezeichnete Farbstoffe sind. Allerdings macht es ihre kurze Bleichzeit beinahe unmöglich, sie für Einzel-Molekül Studien zu gebrauchen.

In Kapitel 3 beschreibe ich ein Experiment, in dem autofluoreszente Proteine zur Untersuchung der Aggregation von  $\beta_2$  adrenergen Rezeptoren ( $\beta_2$ AR) eingesetzt wurden. Diese Rezeptoren sind Teil der en Familie der G-Protein-Gekoppelten Rezeptoren. Der  $\beta_2$ AR war einer der ersten dieser Familie, dessen Sequenz analysiert werden konnte. Daher dient er als ein allgemein akzeptiertes Modellsystem. Der  $\beta_2$ AR spielt eine wichtige Rolle in der Regulierung des Blutdrucks und der Erweiterung der Bronchien. Der Rezeptor kommt in vielen verschieden Geweben und Organen, wie der Niere, Leber und Bauchspeicheldrüse vor. In biochemischen Versuchen werden G-Protein-Gekoppelte Rezeptoren im Allgemeinen in Form von Dimeren gefunden. Geraume Zeit wird ein Zusammenhang zwischen der Dimerisie-

<sup>&</sup>lt;sup>3</sup>Eine Femtosekunde entspricht 0.000'000'000'000'001 Sekunden.

rung und der Funktion vermutet. Es war Teil meiner Forschung, die Dimerisierung eines  $\beta_2$ AR-eYFP Fusions-proteins in abhängigkeit der Stimulierung des Rezeptors zu messen. Dafür habe ich zunächst die Fluoreszenz-Intensitätsverteilung und das Bleichverhalten von eYFP charakterisiert. Die zeitabhängige Verschiebung der Intensitätsverteilung von Rezeptoroligomeren wurde danach modelliert. Wie erwartet, zeigen meine Resulate, dass der hauptsächlich als Dimer vorkommende  $\beta_2$ AR schon wenige Minuten nach der Aktivierung große Aggregate formt. Diese Aggregate bleiben bis zu 10 Minuten lang stabil. Kontrollexperimente mit einem biologisch inaktiven Stereo-Isomer des Liganden des Rezeptors zeigen eine Verschiebung der Intensitätsverteilung hin zu Monomeren, was sich als Effekt des Photobleichens erklären lässt.

Kapitel 4 behandelt die Modifikation unseres Einzel-Molekül Aufbaus für die Detektion von "Rolling Circle Replication" (RCR) DNS Produkten. Mit dem veränderten Aufbau ist es möglich, die DNS-Produkte auch in einem hohen Hintergrund von freien fluoreszenten Labeln in Lösung zu detektieren, ohne dass ein Reinigungsschritt erforderlich ist. RCR ist ein neuartiger DNS-Test, der mit großer Verlässlichkeit die Anwesenheit spezifischer DNS-Sequenzen nachweisen kann. Wenn ein besonderes DNS-"Padlock" ("Vorhangschloss") an seine komplementäre Sequenz bindet, wandelt er sich mit Hilfe einer Ligase zu einer kreisförmigen DNS-Struktur um. Diese kann wiederum durch eine DNS-Polymerase "abgerollt" werden, d.h. es entsteht ein langer Strang von DNS, der aus Kopien der "Padlock"-Sequenz besteht. Mehrere DNS-Label können nun gleichzeitig an dieses RCR-Produkt binden, was die Detektion des DNS-Stranges möglich macht.

Im letzten Kapitel 5 wird eine Methode beschrieben, mit der gleich-



zeitig das Bild und ein Emissionsspektum von fluoreszierenden Objekten aufgenommen werden kann. Dafür wird ein *blazed* Gitter in den Strahlengang des Mikroskops eingeführt. Anders als bei gewöhnlichen Gittern ist es dabei möglich, die Signalstärke in den Beugungsordnungen zu beeinflussen. Wir haben diese neue Methode zunächst an verschieden 20 nm grossen Latexkügelchen getestet, die mit zwei unterschiedlichen Fluoreszenz-Labeln markiert waren. Des Weiteren habe ich diese Methode weiter ausgebaut um Chromosomen mittels sogenannter COBRA-FISH<sup>4</sup> Label eindeutig zu identifizieren.

 $<sup>^4\</sup>mathrm{COmbined}$ Binary RAtio Fluorescence In-Situ Hybridization

## Curriculum Vitae

I was born on November 7, 1973, in the town of Wels, Austria. I grew up in Linz, where I earned my "Matura" at the Bundesrealgymnasium Landwiedstraße in 1992.

After eight months service in the Austrian armed forces, I begun my study of Technical Physics at the Johannes Keppler Universität in Linz in autumn 1993. I was elected student representative in the faculty council for two two-year terms and I also got involved in politics both on the campus as well as the city council as member of the now defunct Liberal Party, where I had to realize that some people conduct political debate in a only marginally less violent way than a full-contact sport. I graduated from the University of Linz with a degree in Physics and the specialization "Biophysics" in November 1999.

During my four year as PhD student in the group of Prof. Thomas Schmidt in Leiden I have had the opportunity to present the results of my research – in the form of a poster – at the 'Frühjahrstagung der Deutschen Physikalischen Gesellschaft, Potsdam, March 2000, the  $3^{rd}$  European Biophysical Congress, München, September 2000, the International Workshop "Single Molecule Detection and Ultra-Sensitive Analysis in Life Science" (Picoquant), Berlin, and the Biophysical Society Annual Meeting. I have been invited to the Symposium of the Daresbury Laboratory, Daresbury, UK, to give a talk "Single Molecule Methods in Biophysical Research", and I am gratefully acknowledging support by the NWO-SMRC Travelgrant 910-31-402, which allowed me to visit the group of Dr. Mats Nilsson at the Rudbeck Laboratories, Uppsala, Sweden.

# Nawoord

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