CIEM5110-2: FEM, lecture 5.2

Introduction to nonlinear material models

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Agenda for today

- 1. Nonlinear material models solver interface
- 2. An overview of different strategies for modeling nonlinear materials
- 3. Illustration: 1D plasticity

CIEM5110-2 workshops and lectures

Recap — Linear FEM

This is the general discretized equilibrium equation:

$$
\underbrace{\int_{\Omega} \mathbf{B}^T \boldsymbol{\sigma} \, d\Omega}_{\mathbf{f}_{int}} = \underbrace{\int_{\Omega} \mathbf{N}^T \mathbf{b} \, d\Omega + \int_{\Gamma_t} \mathbf{N}^T \mathbf{t} \, d\Gamma}_{\mathbf{f}_{ext}}
$$

Assuming linear elasticity, we could substitute $\sigma = \text{DBa}$ to get

$$
\int_{\Omega} \mathbf{B}^T \mathbf{D} \mathbf{B} \, d\Omega \, \mathbf{a} = \int_{\Omega} \mathbf{N}^T \mathbf{b} \, d\Omega + \int_{\Gamma_t} \mathbf{N}^T \mathbf{t} \, d\Gamma \qquad \Rightarrow \qquad \mathbf{K} \mathbf{a} = \mathbf{f}_{\text{ext}}
$$

Linearity is assumed twice there

 $\varepsilon = \mathrm{Ba}$ (kinematic relation)

and

$$
\boldsymbol{\sigma} = \mathbf{D}\boldsymbol{\varepsilon} \quad \text{(constitutive relation)}
$$

Recap — Elasticity

Constant stiffness, full reversibility

In Voigt notation in 2D plane stress:

$$
\mathbf{D} = \begin{bmatrix} \frac{E}{1 - \nu^2} & \frac{E\nu}{1 - \nu^2} & 0\\ \frac{E\nu}{1 - \nu^2} & \frac{E}{1 - \nu^2} & 0\\ 0 & 0 & \frac{E}{2(1 + \nu)} \end{bmatrix}
$$

In Voigt notation in 3D:

Delft

$$
\mathbf{D} = \frac{E}{(1+\nu)(1-2\nu)} \begin{bmatrix} 1-\nu & \nu & \nu & 0 & 0 & 0 \\ \nu & 1-\nu & \nu & 0 & 0 & 0 \\ \nu & \nu & 1-\nu & 0 & 0 & 0 \\ 0 & 0 & 0 & \frac{1-2\nu}{2} & 0 & 0 \\ 0 & 0 & 0 & 0 & \frac{1-2\nu}{2} & 0 \\ 0 & 0 & 0 & 0 & 0 & \frac{1-2\nu}{2} \end{bmatrix}
$$

DAMAGESHR Envelope (max abs) Aug: 75%)

Analysis type: nonlinear analysis (plasticity, damage, contact)

Delft

Damage of the FRP plate

Example: Bolted joints, Fruzsina Csillag (2018)

Objective: investigate the behavior of FRP-steel bolted connections

Deformation of the bolt

Example: Circular micromodels, Pieter Hofman (2021)

Objective: make a micromodel that gives the same failure response in all directions

Analysis type: material nonlinear analysis

Making wind turbines last longer by understanding material behavior:

• Experimental observation: the material loses half of its strength after exposure to hot water

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- Translating small scale phenomena back to the higher scale

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Recap — Nonlinear FEM

Discretized form:

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\int_{\Omega} \mathbf{B}^T \boldsymbol{\sigma} \, d\Omega = \int_{\Omega} \mathbf{N}^T \mathbf{b} \, d\Omega + \int_{\Gamma_t} \mathbf{N}^T \mathbf{t} \, d\Gamma
$$

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Geometric nonlinearity:

 $\mathbf{B} = \mathbf{B} (\mathbf{a})$

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$$

Geometric nonlinearity:

 $\mathbf{B} = \mathbf{B} (\mathbf{a})$

Material nonlinearity:

$$
\frac{\partial \pmb{\sigma}}{\partial \pmb{\varepsilon}} = \mathbf{D} = \mathbf{D}\left(\mathbf{a}\right)
$$

Load control with a nonlinear material

Require: Nonlinear relation $\mathbf{f}_{\mathrm{int}}(\mathbf{a})$ with $\mathbf{K}(\mathbf{a}) = \frac{\partial \mathbf{f}_{\mathrm{int}}}{\partial \mathbf{a}}$

1: Initialize $n \leftarrow 0$, $\mathbf{a}^0 \leftarrow \mathbf{0}$

- 2: **while** n < number of time steps **do**
- 3: Get new external force vector: $\mathbf{f}^{n+1}_{\text{ext}}$ ext
- 4: Initialize new solution at old one: $\mathbf{a}^{n+1} \leftarrow \mathbf{a}^n$
- 5: Compute internal force and stiffness: ${\bf f}^{n+1}_{\rm int}({\bf a}^{n+1})$, ${\bf K}^{n+1}({\bf a}^{n+1})$
- 6: Evaluate first residual: $\mathbf{r} = \mathbf{f}_{\text{ext}}^{n+1} \mathbf{f}_{\text{int}}^{n+1}$ int

7: **repeat**

- 8: Solve linear system of equations: $\mathbf{K}^{n+1}\Delta \mathbf{a} = \mathbf{r}$
- 9: Update solution: $\mathbf{a}^{n+1} \leftarrow \mathbf{a}^{n+1} + \Delta \mathbf{a}$
- 10: Compute internal force and stiffness: ${\bf f}^{n+1}_{\rm int}({\bf a}^{n+1})$, ${\bf K}^{n+1}({\bf a}^{n+1})$
- 11: Evaluate residual: $\mathbf{r} = \mathbf{f}_{\text{ext}}^{n+1} \mathbf{f}_{\text{int}}^{n+1}$ int
- 12: **until** |r| < tolerance
- 13: $n \leftarrow n + 1$

14: **end while**

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elft

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$$
\begin{aligned}\n & \mathbf{a} \\
 & \downarrow \\
 & \varepsilon = \mathbf{B}\mathbf{a} \\
 & \downarrow \\
 & \sigma = \sigma(\varepsilon) \mathbf{D} = \frac{\partial \sigma}{\partial \varepsilon} \\
 & \downarrow \\
 & \mathbf{f}_{\text{int}} = \int_{\Omega} \mathbf{B}^{\text{T}} \sigma \mathrm{d}\Omega \quad \mathbf{K} = \int_{\Omega} \mathbf{B}^{\text{T}} \mathbf{D} \mathbf{B} \mathrm{d}\Omega\n \end{aligned}
$$

Hyperelasticity

Analogous to a nonlinear spring

Popular for modeling large strains

Still fully reversible

Usually derived from a single scalar potential W

Viscoelasticity

Stiffness is time-dependent

Fully reversible response, but stiffer if loaded faster

An integral in time appears

Damage

Loss of load-carrying area modeled as loss of stiffness

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Loss of load-carrying area modeled as loss of stiffness

$$
A = (1 - d) A_0
$$

The damage d evolves according to a loading function:

$$
f(\tilde{\varepsilon}, \alpha) = \tilde{\varepsilon} - \alpha, \quad f \le 0, \dot{\alpha} \ge 0, f\dot{\alpha} = 0
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Loss of load-carrying area modeled as loss of stiffness

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The damage d evolves according to a loading function:

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f(\tilde{\varepsilon}, \alpha) = \tilde{\varepsilon} - \alpha, \quad f \le 0, \dot{\alpha} \ge 0, f\dot{\alpha} = 0
$$

and an evolution equation:

 $d = d(\alpha)$

Displacement control with a history-dependent material

Require: Nonlinear relation $f_{\text{int}}(a)$ with $K(a) = \frac{\partial f_{\text{int}}}{\partial a}$

- 1: Initialize new solution at old one: $\mathbf{a}^{n+1} \leftarrow \mathbf{a}^n$
- 2: Update material model: $\{\bm \sigma^{n+1}, \mathbf D^{n+1}, \bm \alpha_\text{new}\} = \mathcal{M}\left(\bm \varepsilon^{n+1}, \bm \alpha_\text{old}\right)$
- 3: Compute internal force and stiffness: $\mathbf{f}^{n+1}_{\rm int} =$ z Ω $\mathbf{B}^{\mathrm{T}} \boldsymbol{\sigma}^{n+1} d\Omega$, $\mathbf{K}^{n+1} = \mathbf{A}$ Ω $\mathbf{B}^{\mathrm{T}}\mathbf{D}^{n+1}\mathbf{B}\mathrm{d}\Omega$
- 4: Constrain ${\bf K}^{n+1}$ so that $\Delta {\bf a}_c = \overline{\bf a}^{n+1} \overline{\bf a}^n$
- 5: Evaluate first residual: $\mathbf{r} = -\mathbf{f}^{n+1}_{\text{int},f}$ int,f

6: **repeat**

- 7: Solve linear system of equations: $\mathbf{K}^{n+1}\Delta\mathbf{a} = \mathbf{r}$
- 8: Update solution: $\mathbf{a}^{n+1} \leftarrow \mathbf{a}^{n+1} + \Delta \mathbf{a}$
- 9: Update material model: $\{\boldsymbol{\sigma}^{n+1},\mathbf{D}^{n+1},\boldsymbol{\alpha}_{\text{new}}\} = \mathcal{M}\left(\boldsymbol{\varepsilon}^{n+1},\boldsymbol{\alpha}_{\text{old}}\right)$
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- 11: Evaluate residual: $\mathbf{r} = -\mathbf{f}^{n+1}_{\text{int},f}$ int,f
- 12: Constrain \mathbf{K}^{n+1} so that $\Delta \mathbf{a}_c = 0$
- 13: **until** |r| < tolerance

14: Commit material history: $\alpha_{old} \leftarrow \alpha_{new}$ $\mathcal{A}_{\mathcal{C}}$ **Delft**

Discontinuous damage

Model an actual displacement discontinuity

Traction-separation through a cohesive zone law

Special interface elements are needed

Similar to damage, but explicit link to energy dissipation

$$
\begin{array}{|c|c|c|}\hline \text{non-odd} & \text{non-odd} \\ \hline \text{non-odd} & \text{non-odd} \\ \hline \text{non-odd} & \text{non-odd} \\ \hline \end{array}
$$

$$
\begin{array}{ccc}\n\left[\mathbf{u}\right] = \mathbf{u}^{\text{top}} - \mathbf{u}^{\text{bot}} & \Rightarrow & \left[\mathbf{N} - \mathbf{N}\right] \begin{bmatrix} \mathbf{a}^{\text{top}} \\ \mathbf{a}^{\text{bot}} \end{bmatrix} & \Rightarrow & \left[\mathbf{u}\right] = \mathbf{N}_{\Gamma} \mathbf{a}_{\Gamma} \\
\downarrow & & \\
\hline\n\end{array}
$$
\nation

\n
$$
\mathbf{t} = \mathcal{T}\left(\begin{bmatrix} \mathbf{u} \end{bmatrix}\right) & \Rightarrow & \mathbf{f}_{\text{int}}^{\Gamma} = \int_{\Gamma} \mathbf{N}_{\Gamma}^{\text{T}} \mathbf{t} \mathrm{d}\Gamma
$$

Plasticity

Deformations are split:

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\boldsymbol{\varepsilon}=\boldsymbol{\varepsilon}^\mathrm{e}+\boldsymbol{\varepsilon}^\mathrm{p}
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 $f\left(\boldsymbol{\sigma}\right) = \tilde{\boldsymbol{\sigma}} - \sigma_{y}\left(\varepsilon_{\rm acc}^{\rm p}\right)$

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A yield surface defines the plasticity threshold:

 $f\left(\boldsymbol{\sigma}\right) = \tilde{\boldsymbol{\sigma}} - \sigma_{y}\left(\varepsilon_{\rm acc}^{\rm p}\right)$

Plastic flow occurs when the yield surface is pushed:

 $f=0, \dot{f}=0 \quad \Rightarrow \quad \dot{\boldsymbol{\varepsilon}}^{\mathrm{p}} = \dot{\gamma} \mathbf{m}$

We need to keep track of history variables $\boldsymbol{\alpha} = [\boldsymbol{\varepsilon}^\mathrm{p}, \varepsilon_{\mathrm{acc}}^\mathrm{p}]$

Hardening plasticity for bending moments

Curvatures are split:

 $\kappa = \kappa^{\mathrm{e}} + \kappa^{\mathrm{p}}$

A yield point defines the plasticity threshold:

$$
f(M, \kappa^{\rm p}) = |M| - M^{\rm y}(\kappa^{\rm p}_{\rm acc}), \quad \kappa^{\rm p}_{\rm acc} = \int_t |\dot{\kappa}^{\rm p}| {\rm d}t
$$

Plastic flow occurs when the yield surface is pushed:

 $f = 0, \dot{f} = 0 \Rightarrow \Delta \kappa^{\rm p} = \Delta \gamma \operatorname{sign}(M)$

History variables are $\boldsymbol{\alpha} = [\kappa^{\mathrm{p}}, \kappa^{\mathrm{p}}_{\mathrm{acc}}]$

 $M = EI(\kappa - \kappa^{\mathbf{p}})$

 $\widetilde{\mathbf{T}}$ U **Delft** Downwards displacement

Recap and outlook

This short story on nonlinear FEM is now round:

- Sources of nonlinearity
- Dealing with geometric nonlinearities
- Dealing with material nonlinearities
- Solver strategies for nonlinear FEM

If you want to keep digging, this is just the beginning:

- CIEM5210-2: Advanced constitutive modeling
	- Computational plasticity, continuum and discrete damage, advanced path following methods
- CIEM1303 (upscaling techniques) and CIEM1301 (advanced computational mechanics):
	- Multiscale/multiphysics modeling, recent advances in nonlinear FEM

